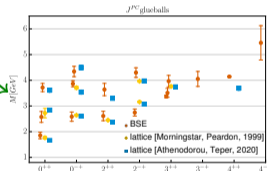
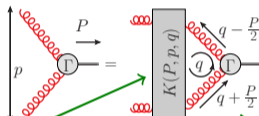
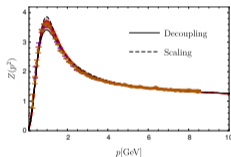
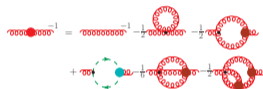
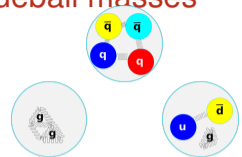


On the role of truncations of the 3PI effective action for glueball masses

Markus Q. Huber

Institute of Theoretical Physics, Giessen University



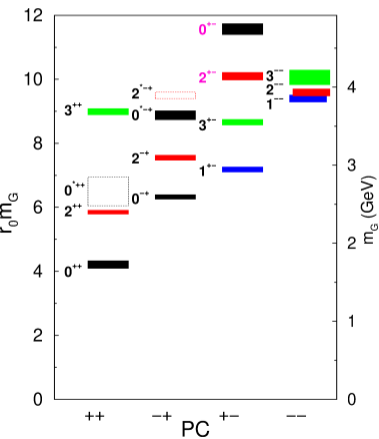
QuantFunc

September 4, 2024

Valencia, Spain

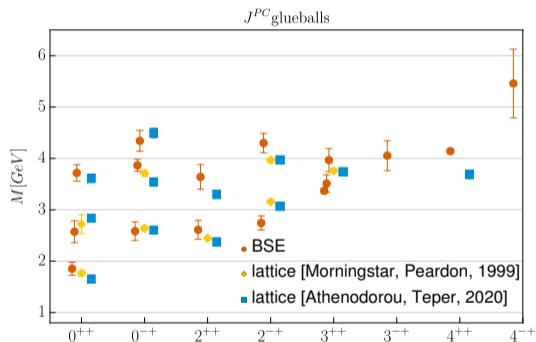
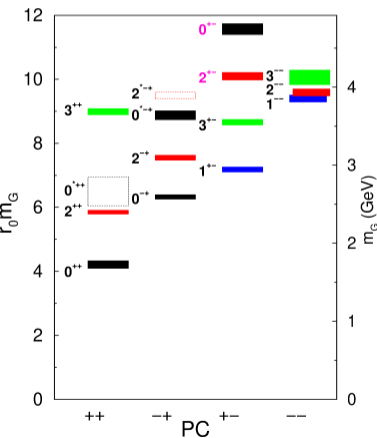
Glueball results

Pure gauge theory glueballs known [Morningstar, Peardon, Phys. Rev. D60 (1999)].



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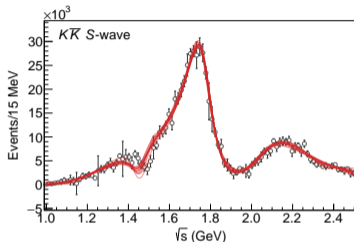
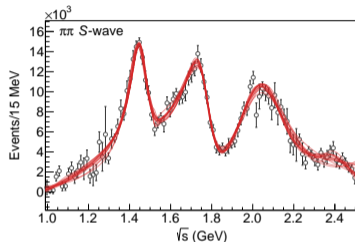


Lattice update: [Athenodorou, Teper, JHEP11 (2020)]

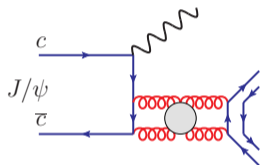
Functional results: [MQH, Fischer, Sanchis-Alepuz, Eur.Phys.J.C81 (2021)]

→ Agreement with lattice results.

Glueballs from J/ψ decay



[JPAC Coll., Rodas et al.,
Eur.Phys.J.C 82 (2022)]



Scalar glueball candidate: Coupled-channel analyses of exp. data (BESIII)

- +add. data, largest overlap with $f_0(1770)$
- largest overlap with $f_0(1710)$

[Sarantsev, Denisenko, Thoma, Klempt, Phys. Lett. B 816 (2021)]

[JPAC Coll., Rodas et al., Eur.Phys.J.C 82 (2022)]

Pseudoscalar glueball candidate:

- $X(2370)$

[Ablikim et al. (BESIII), PRL132 (2024)]

Challenges/Reliability

Lattice:

- Statistical method \rightarrow error bars
- Continuum extrapolation

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- Time-like momenta \rightarrow Extrapolation errors

Challenges/Reliability

Yang-Mills:

Lattice:

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- Continuum extrapolation

Functional:

- Truncation errors
- Time-like momenta \rightarrow Extrapolation errors

QCD:

Lattice:

- Statistical method, but poor signal-to-noise ratio
- Continuum extrapolation
- Operator basis
- Mixing between operators
- Physical quark masses

Functional:

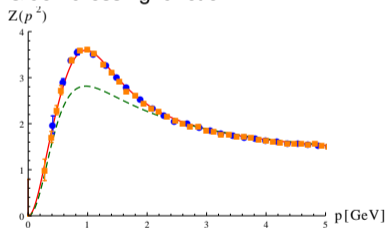
- Truncation errors
- Time-like momenta \rightarrow Extrapolation errors
- Mixing between operators

Input for bound state equations: 2013?

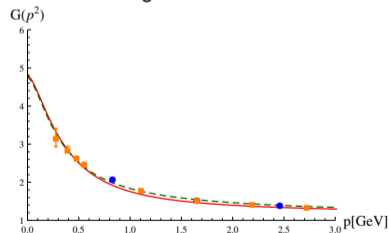
Bound state equations require propagators and vertices as input.

Propagators from 2013 [Sternbeck, hep-lat/0609016; MQH, von Smekal, JHEP 04 (2013)]:

Gluon dressing function:



Ghost dressing function:

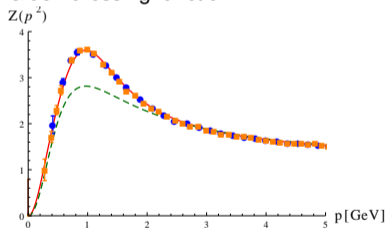


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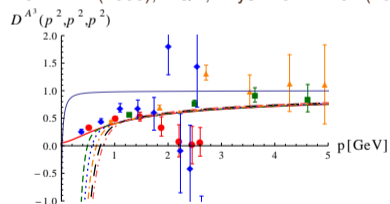
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Gluon dressing function:



Three-gluon vertex [Cucchieri, Maas, Mendes, Phys. Rev. D 77 (2008); MQH, Phys. Rev. D 101 (2020)]:



Used a three-gluon vertex model tailored to reproduce lattice results.

→ Insufficient for bound states.

Model based BSE calculations ($J = 0$)

- [Meyers, Swanson, Phys.Rev.D87 (2013)]
- [Sanchis-Alepuz, Fischer, Kellermann, von Smekal, Phys.Rev.D92, (2015)]
- [Souza et al., Eur.Phys.J.A56 (2020)]
- [Kaptari, Kämpfer, Few Body Syst.61 (2020)]

→ Calculations possible, but no quantitative prediction yet.

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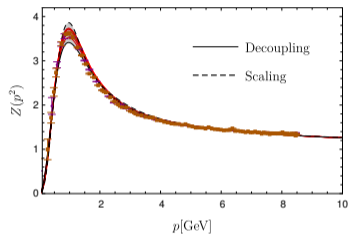
→ Calculations possible, but no quantitative prediction yet.

Better models? Alternative: Self-consistently calculated input.

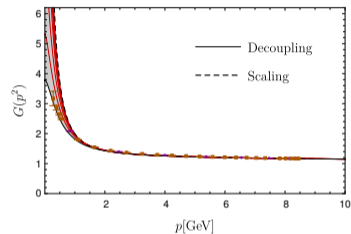
Input for bound state equations: 2020?

[MQH, Phys. Rev. D 101 (2020)]

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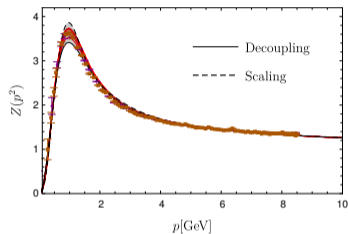
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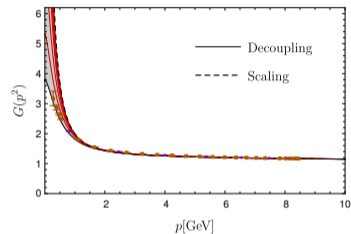
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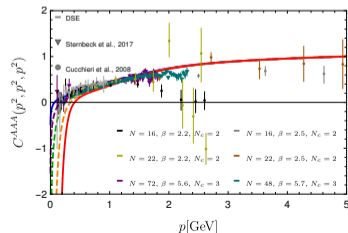
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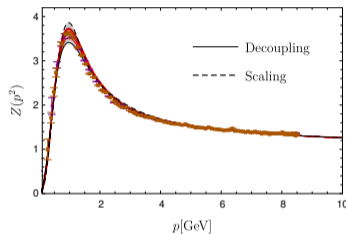
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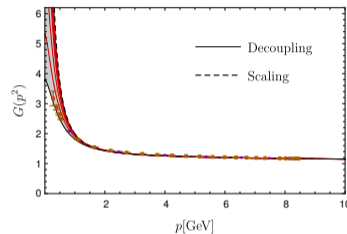
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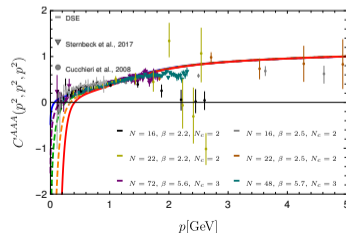
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Three-gluon vertex [Cucchieri, Maas, Mendes, Phys. Rev. D 77 (2008);
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Still truncations in equations. . .

How far should we trust the results?



Overview

How to assess the quality of a truncation?

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- Agreement with analytic results (perturbation theory)?

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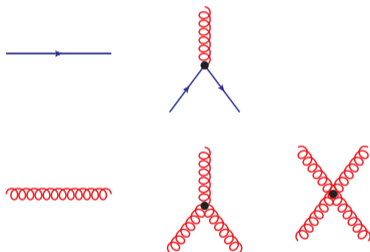
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- 3PI effective action and its truncation
- Correlation functions: equations of motion, truncations, impact of extensions
- Bound state equations: kernels and their truncations, impact of extensions

Effective actions

QCD: non-Abelian gauge theory



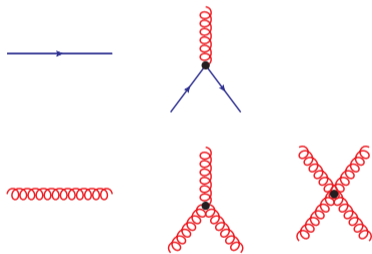
Lagrangian density of QCD \mathcal{L}_{QCD} :

→ 1PI effective action $\Gamma[\bar{q}, q, A, \bar{c}, c]$:

- Generating functional of correlation functions
- Equations of motion: DSEs

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Lagrangian density of QCD \mathcal{L}_{QCD} :

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→ 3PI effective action $\Gamma[\bar{q}, q, A, \bar{c}, c, D, \Gamma^{(3)}]$:

- Equations of motion from stationarity conditions
- Nonperturbative loop-expansion

→ three-loop truncation: $\Gamma^{3l} = \Gamma^{0,3l} + \Gamma^{\text{int},3l}$

$$\Gamma^{0,3l}[\Phi, D, \Gamma^{(3)}] = \frac{1}{8} \text{ (two loops)} + \frac{1}{6} \text{ (one loop)} - \frac{1}{4} \text{ (one loop)} + \frac{1}{8} \text{ (triangle loop)}$$

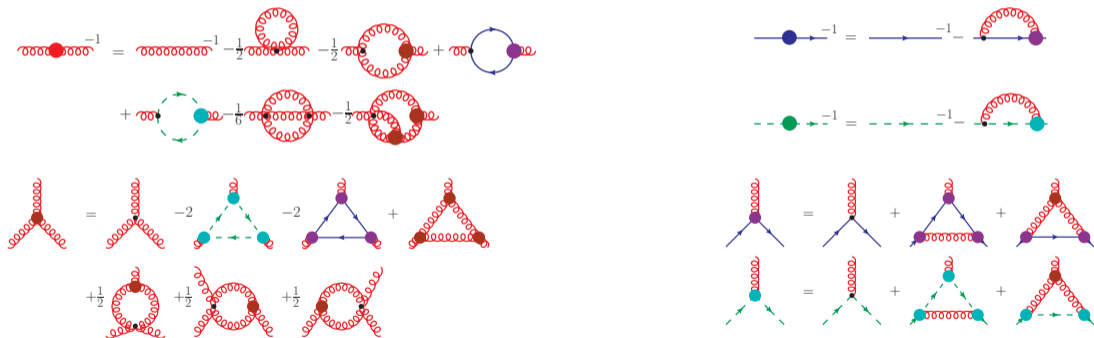
$$\Gamma^{\text{int},3l}[D, \Gamma^{(3)}] = -\frac{1}{12} \text{ (triangle loop)} + \frac{1}{2} \text{ (triangle loop)} + \frac{1}{24} \text{ (triangle loop)} - \frac{1}{3} \text{ (triangle loop)} - \frac{1}{4} \text{ (triangle loop)}$$

The equations show the decomposition of the three-loop truncation of the 3PI effective action into tree-level and loop-level diagrams. The first equation, $\Gamma^{0,3l}$, includes a two-loop diagram (two circles), a one-loop diagram (one circle with a horizontal line), another one-loop diagram (one circle with a vertical line), and a triangle loop diagram. The second equation, $\Gamma^{\text{int},3l}$, includes several diagrams with dashed green lines representing internal three-particle irreducible (3PI) vertices, including triangle loops and more complex topologies.

[Berges, Phys. Rev. D 70 (2004); Carrington, Gao, Phys. Rev. D 83 (2011)]

Equations of motion

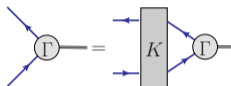
Stationarity conditions: $\frac{\delta\Gamma}{\delta D} = 0$, $\frac{\delta\Gamma}{\delta\Gamma^{(3)}} = 0 \rightarrow$ Equations of motion (already truncated):



(Similar to DSEs, but still different.)

Bound state equations

Generic form:

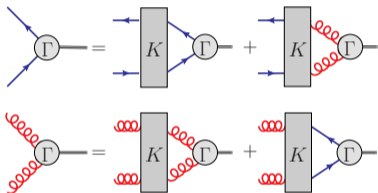


Kernels can be obtained from the 3PI effective action [Fukuda, Prog. Theor. Phys 78 (1987); McKay, Munczek, Phys. Rev. D 40 (1989); Sanchis-Alepuz, Williams, J. Phys: Conf. Ser. 631 (2015)]:

$$K = -2 \frac{\delta^2 \Gamma^3}{\delta D^2}$$

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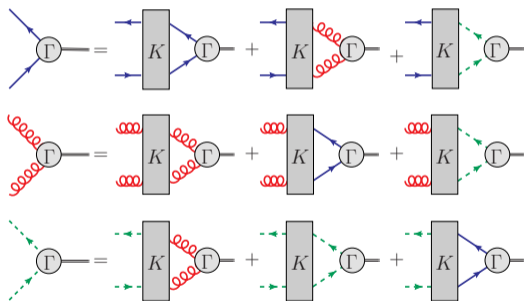


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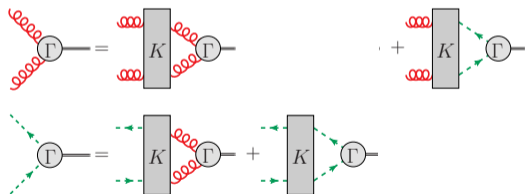
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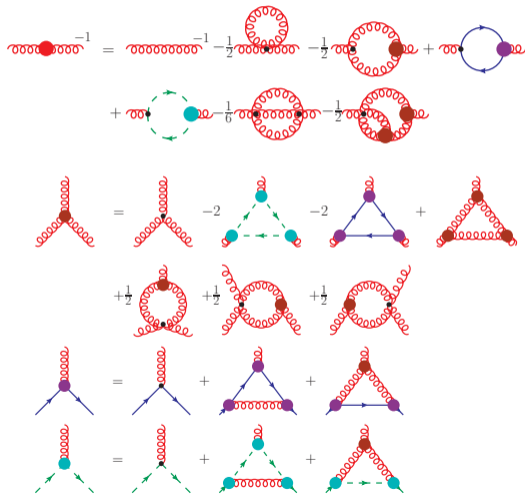
Focus on pure glueballs.



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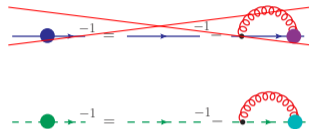
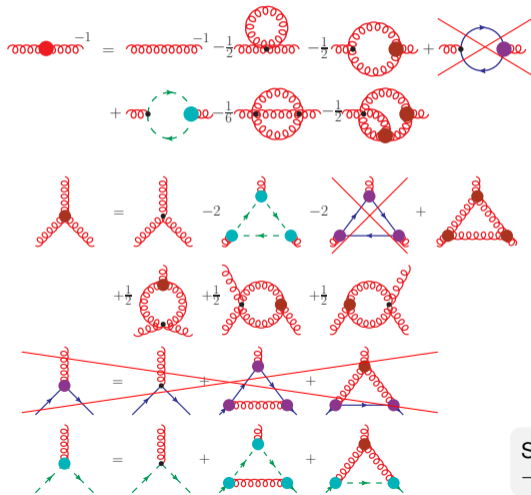
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Equations of motion from 3PI–3-loop



- Truncation at level of action \rightarrow System of equations complete!
- \rightarrow Self-contained: Only parameters are the **strong coupling and the quark masses!**
- Parts: propagators, ghost-gluon vertex, three-gluon vertex, four-gluon vertex
 - Individual calculations
 - Combination with lattice results
 - Complete setup

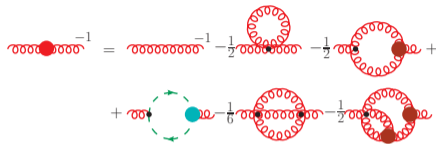
Equations of motion from 3PI-3-loop



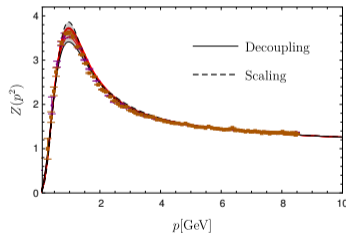
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Start with **pure gauge theory**.
 \rightarrow [MQH, Phys.Rev.D 101 (2020)]

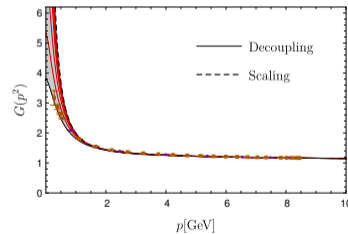
Landau gauge propagators



Gluon dressing function:



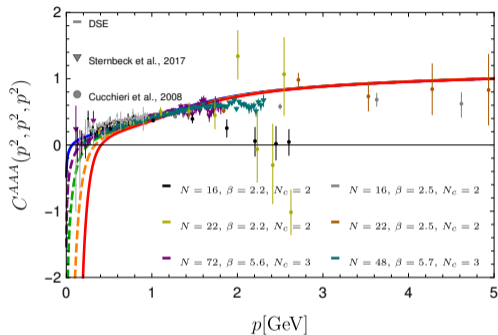
Ghost dressing function:



[Sternbeck, hep-lat/0609016; MQH, Phys. Rev. D 101 (2020)]

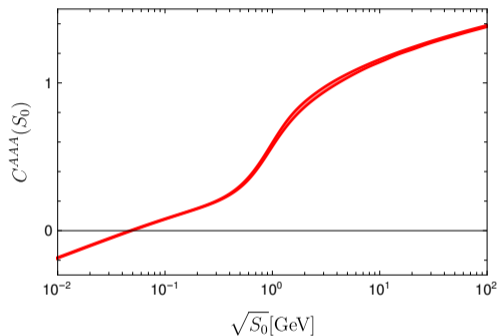
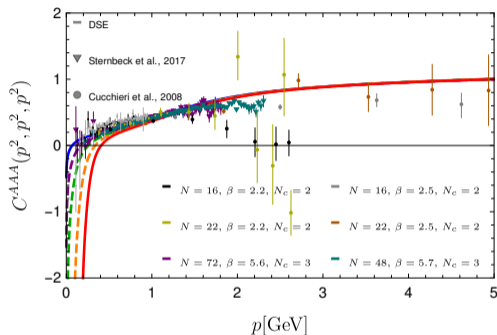
(Family of solutions [von Smekal, Alkofer, Hauck, PRL79 (1997); Aguilar, Binosi, Papavassiliou, Phys.Rev.D 78 (2008); Boucaud et al., JHEP06 (2008); Fischer, Maas, Pawłowski, Ann.Phys. 324 (2008); Alkofer, MQH, Schwenzer, Phys. Rev. D 81 (2010)])

Three-gluon vertex: Kinematics



- IR suppression with zero crossing

Three-gluon vertex: Kinematics



- IR suppression with zero crossing
- Simple kinematic dependence (singlet variable S_0 of S_3), “planar degeneracy”

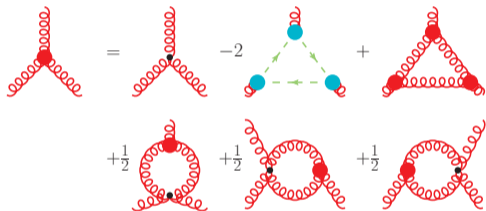
First observation: [Eichmann, Williams, Alkofer, Vujanovic, Phys.Rev.D89 (2014)], but already in old data [Blum, Huber, Mitter, von Smekal, Phys.Rev.D89 (2014)]; lattice: [Pinto-Gómez et al., Phys.Lett.B838 (2023)]; [Aguilar et al., Eur.Phys.J.C83 (2023)]

→ Talks by de Soto, Rodríguez-Quintero

[Cucchieri, Maas, Mendes, Phys. Rev. D
77 (2008); Sternbeck et al., 1702.00612;
MQH, Phys. Rev. D 101 (2020)]

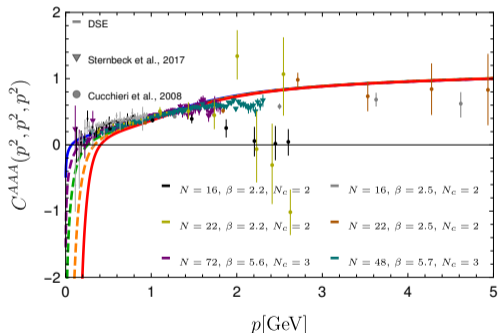
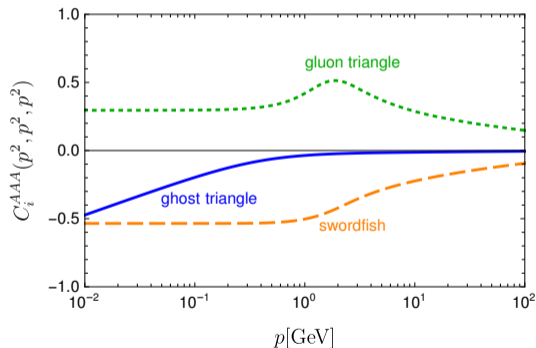
Three-gluon vertex: Individual diagrams

Importance of individual diagrams?



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Importance of individual diagrams?



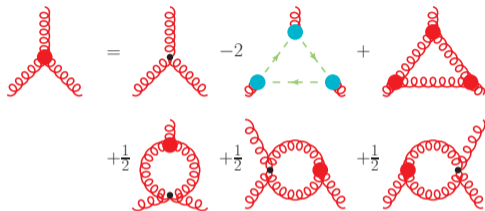
[MQH, Phys. Rev. D 101 (2020)]

→ Cancellations between diagrams important.

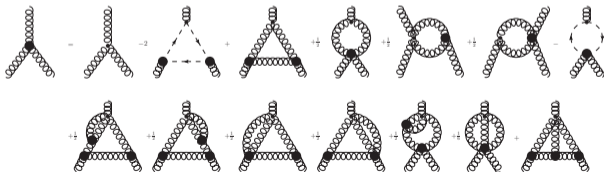
Three-gluon vertex: 1PI vs. 3PI

[MQH, Phys. Rev. D 101 (2020)]

3PI:

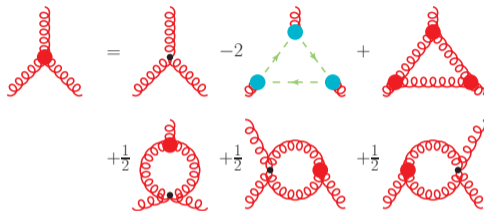


1PI:

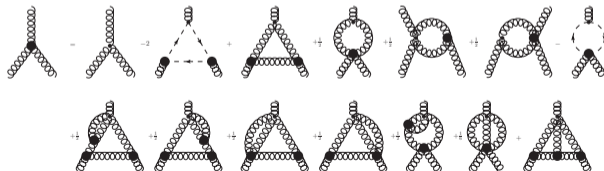


Three-gluon vertex: 1PI vs. 3PI

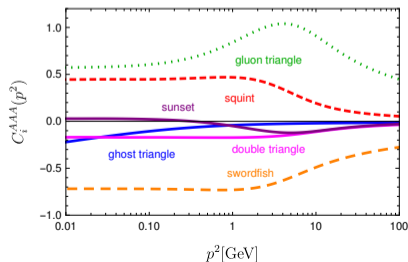
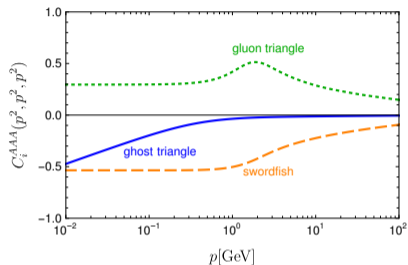
3PI:



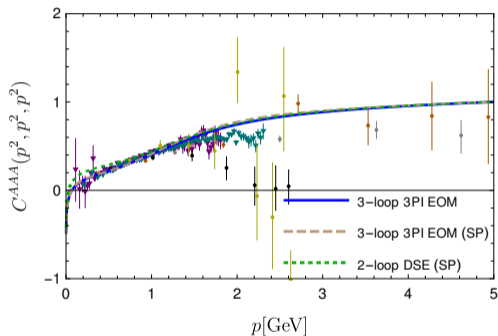
1PI:



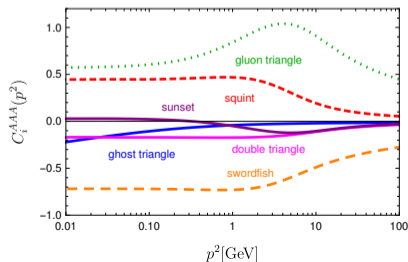
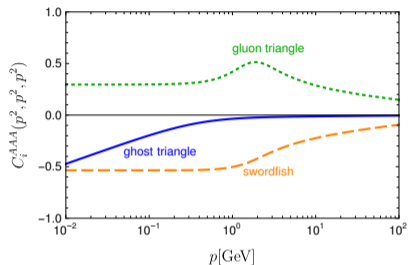
[MQH, Phys. Rev. D 101 (2020)]



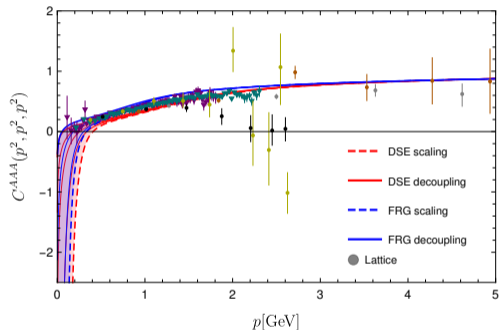
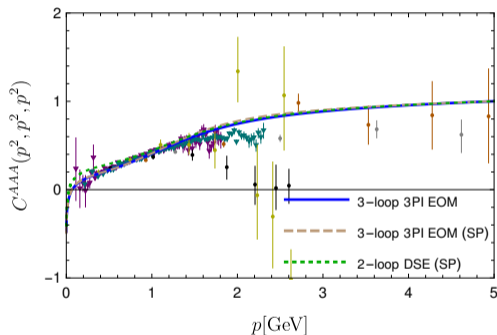
Three-gluon vertex: 1PI vs. 3PI



[MQH, Phys. Rev. D 101 (2020)]



Three-gluon vertex: More methods



[Cucchieri, Maas, Mendes, Phys.Rev.D77 (2008); Sternbeck et al., Proc.Sci. LATTICE2016 (2017); Cyrol et al., Phys.Rev.D 94 (2016); MQH, Phys.Rev.D101 (2020)]

→ Agreement between lattice, FRG, DSE and 3PI.

Stability of the three-gluon vertex results

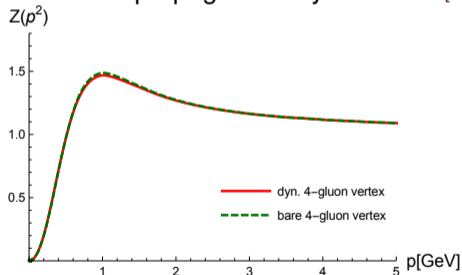
- Agreement with lattice and other functional results. ✓

Stability of the three-gluon vertex results

- Agreement with lattice and other functional results. ✓
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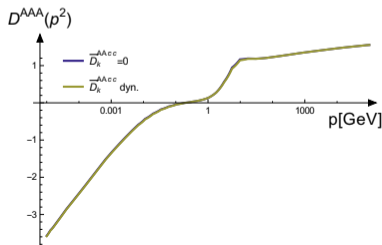
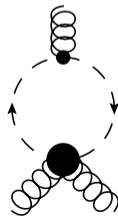
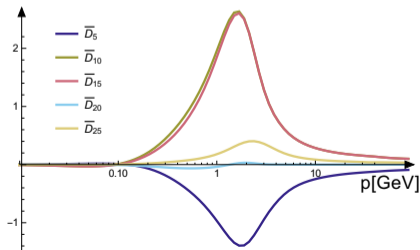
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- Two-ghost-two-gluon vertex with 25 dressings [MQH, Eur. Phys.J.C77 (2017)]: ✓
(FRG: [Corell, SciPost Phys. 5 (2018)])



Four-gluon vertex

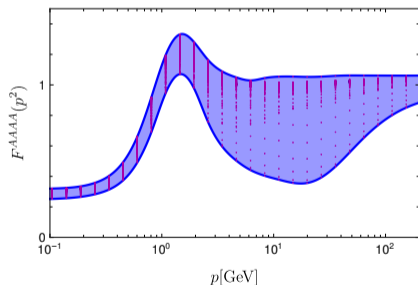
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Four-gluon vertex from DSE using 3PI input:

[MQH, Phys. Rev. D 101 (2020)]



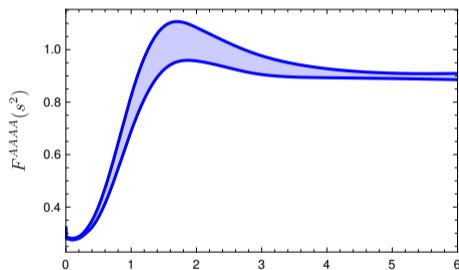
Calculated for three kinematic variables (singlet and doublet of S_4).

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$$s = \sqrt{\frac{3}{2}} S_0 \quad [\text{GeV}]$$

Shown: Subset of one vanishing gluon momentum.

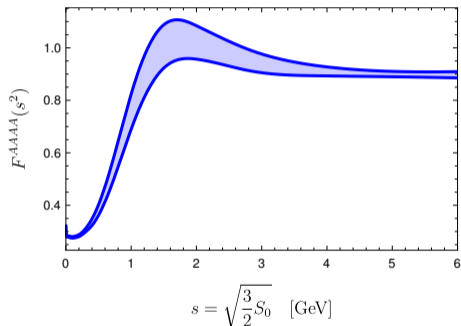
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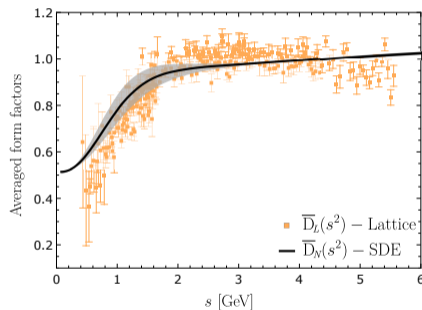
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Talks by Oliveira, de Soto, Santos [Colaco, Oliveira Silva, Phys. Rev. D 109 (2024); Aguilar et al., 2408.06135]

lattice, 4PI effective action → possibilities for checks between three methods

Possible approximations for bound state equations

Basic and complexity considerations:

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Tensor basis:

Full? Relevant subset?

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Specify diagrams based on. . .

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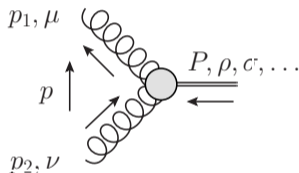
Input:

- Availability
- Kinematics
- Tensors

Tensor bases for glueballs

[MQH, Fischer, Sanchis-Alepuz, Eur.Phys.J.C81 (2021)]

$$\Gamma_{\mu\nu\rho\sigma\dots}(p_1, p_2) = \sum \tau_{\mu\nu\rho\sigma\dots}^i(p_1, p_2) h_i(p_1, p_2)$$



Increase in complexity:

- 2 gluon indices (transverse)
- J spin indices (symmetric, traceless, transverse to P)
- Linear independence (nontrivial for $P = -1$)

Numbers of tensors:

J	$P = +1$	$P = -1$
0	2	1
1	4	3
>2	5	4

Low number of tensors, but high-dimensional tensors!

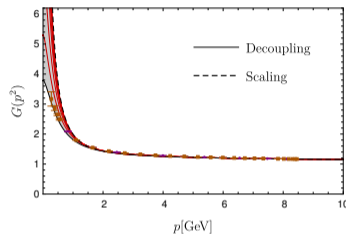
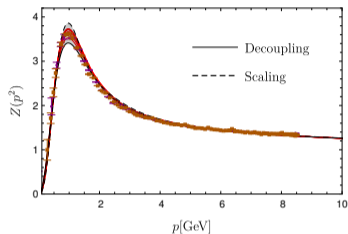
→ Computational cost increases with J .

→ Full tensor basis used.

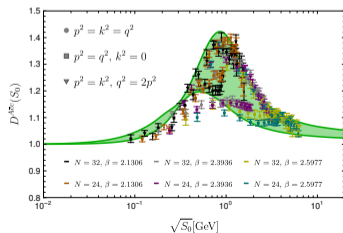
Input

Propagators: Full

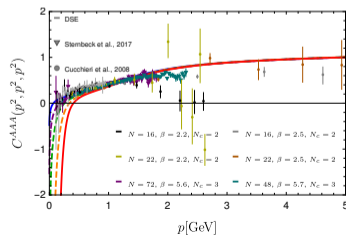
[MQH, Phys. Rev. D 101 (2020)]



Ghost-gluon vertex: Full



Three-gluon vertex: Leading tensor, full kinematics



Kernels from 3PI–3-loop

One-particle exchange kernels, e.g., ladder truncation (long-time work horse):



Kernels from 3PI-3-loop

One-particle exchange kernels, e.g., ladder truncation (long-time work horse):



$$K = \text{Diagram 1} + \frac{1}{2} \text{Diagram 2} + \frac{1}{2} \text{Diagram 3} - \text{Diagram 4} + \text{Diagram 5} + \text{Diagram 6} + \frac{1}{2} \text{Diagram 7}$$

$$K = \text{Diagram 8} + \frac{1}{2} \text{Diagram 9} + \frac{1}{2} \text{Diagram 10}$$

$$K = \text{Diagram 11} + \frac{1}{2} \text{Diagram 12}$$

$$K = \text{Diagram 13} + \frac{1}{2} \text{Diagram 14} + \frac{1}{2} \text{Diagram 15}$$

→ Two-loop diagrams in BSEs!

Results from one-particle exchange kernels

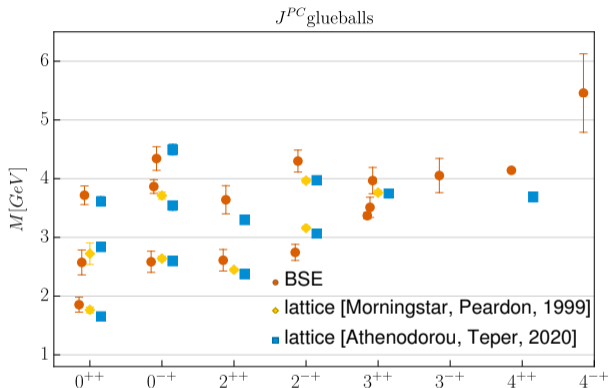
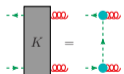
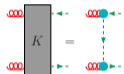
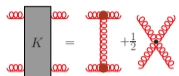
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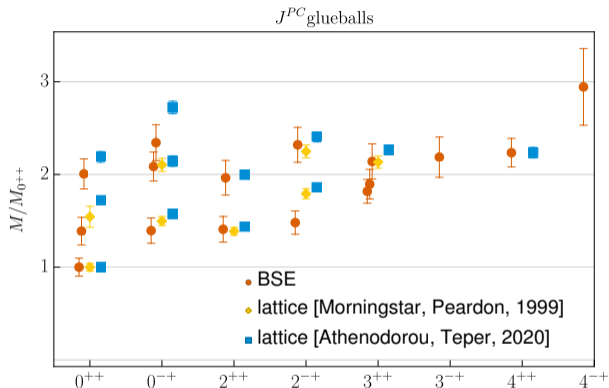
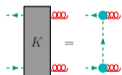
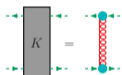
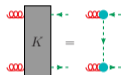
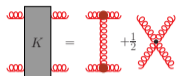
Results from one-particle exchange kernels



[MQH, Fischer, Sanchis-Alepuz, Eur.Phys.J.C81 (2021)]

Spectral reconstruction [Pawlowski et al., Phys.Rev.D 108 (2023)]:
 0^{++} : 1870 MeV, 0^{-+} : 2700 MeV

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Beyond one-particle exchange

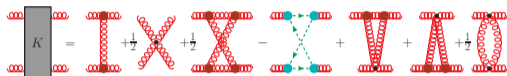
Gluon-gluon interactions leading \rightarrow Test its full 3PI–3-loop kernel:

$$K = \text{Tree} + \frac{1}{2} \text{Loop} + \frac{1}{2} \text{TwoLoop} - \text{ThreeLoop} + \text{TwoLoop} + \text{OneLoop}$$

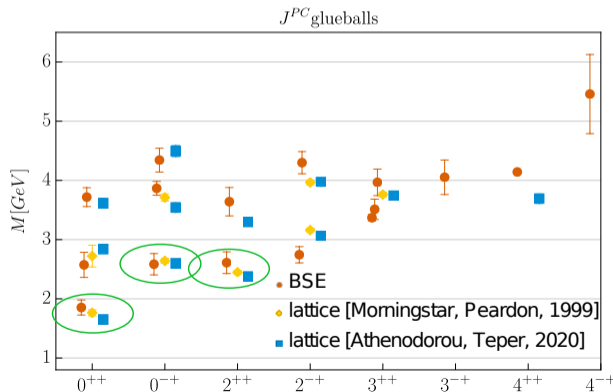
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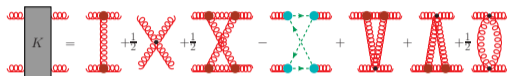


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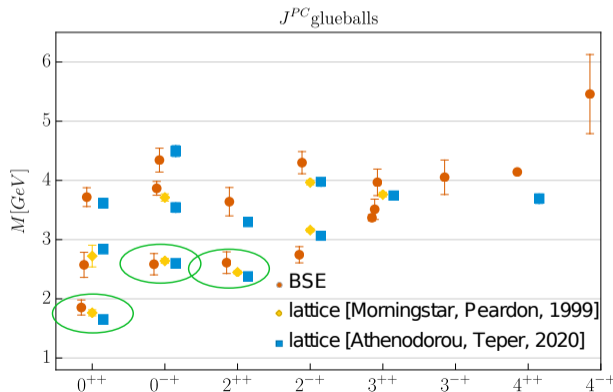


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● 0^{-+} : none

[MQH, Fischer, Sanchis-Alepuz, EPJ Web Conf. 258 (2022)]

● 0^{++} : $< 2\%$

[MQH, Fischer, Sanchis-Alepuz, HADRON2021, Rev.Mex.Fis.Suppl. 3 (2022)]

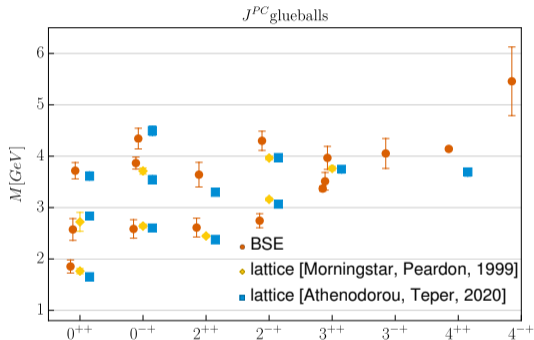
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Summary

Glueballs from functional equations

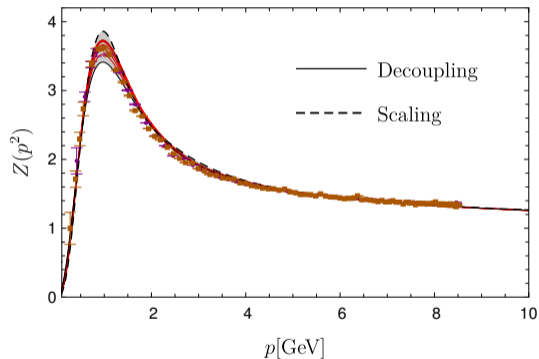
- Tool for hadron physics: From qualitative insights to quantitative results



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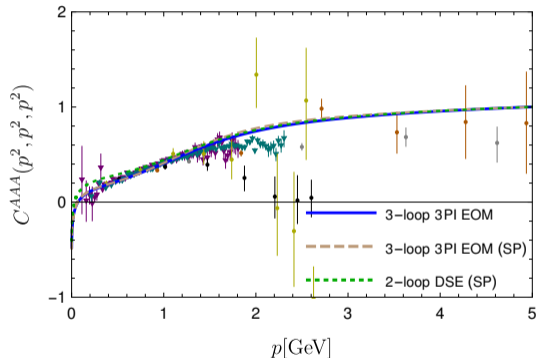
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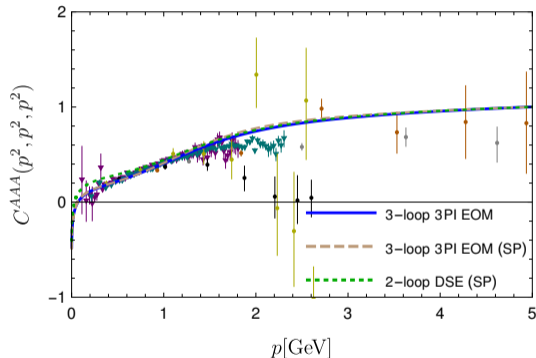
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Diagram illustrating the kernel K as a sum of two diagrams: a vertical line and a cross, with a coefficient of $+\frac{1}{2}$ for the cross term.

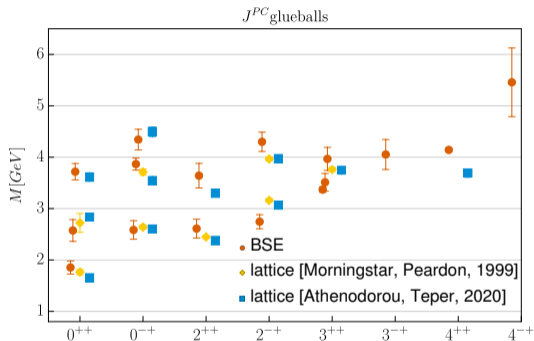
vs.

Diagram illustrating the kernel K as a sum of six diagrams: a vertical line, a cross, a double cross, a diamond, a V-shape, and a loop, with coefficients of $+\frac{1}{2}$, $+\frac{1}{2}$, $-$, $+$, and $+\frac{1}{2}$ respectively.

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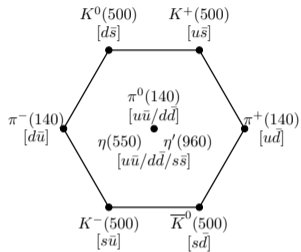


Thank you for your attention!

Scalar sector

$J^{PC} = 0^{++} \rightarrow q\bar{q}$ mesons, tetraquarks and glueballs

$$m_u \sim m_d < m_s$$

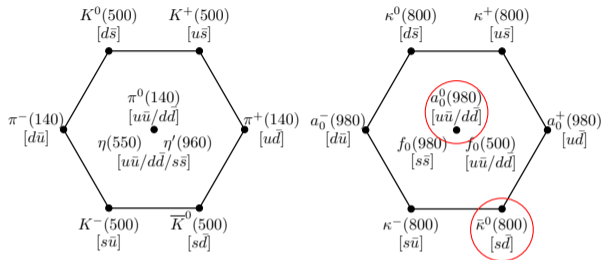


Pseudoscalar:
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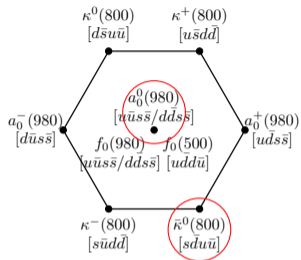
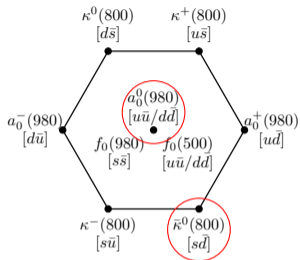
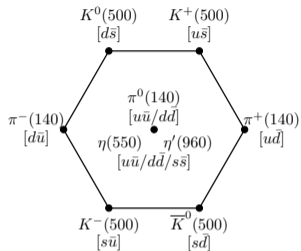
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Functional review:

[Eichmann, Fischer,
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Few-Body Syst.61 (2020)]

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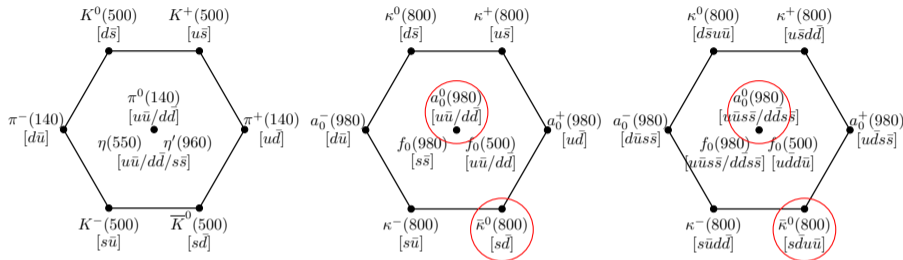
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Inverted mass hierarchy if tetraquarks
[Jaffe, Phys. Rev. D 15 (1977)]

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[Jaffe, Phys. Rev. D 15 (1977)]

Scalar glueballs: $Q = 0$, isoscalar

$f_0(500)$	$f_0(1370)$
$f_0(980)$	$f_0(1500)$
	$f_0(1710)$

glueball candidates

$J = 1$ glueballs

Landau-Yang theorem

Two-photon states cannot couple to $J^P = 1^\pm$ or $(2n + 1)^-$

[Landau, Dokl.Akad.Nauk SSSR 60 (1948); Yang, Phys. Rev. 77 (1950)].

(\rightarrow Exclusion of $J = 1$ for Higgs because of $h \rightarrow \gamma\gamma$.)

Applicable to glueballs?

\rightarrow Not in this framework, since gluons are not on-shell.

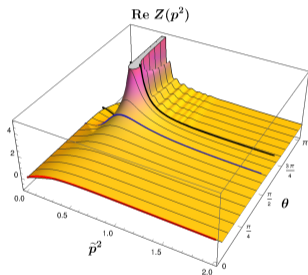
\rightarrow Presence of $J = 1$ states is a dynamical question.

$J = 1$ not found here.

Correlation functions in the complex plane

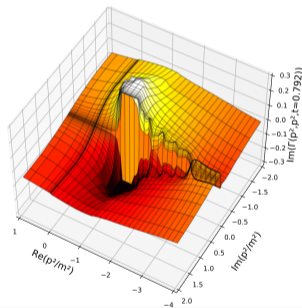
Gluon

[Fischer, MQH, Phys.Rev.D 102 (2020)]



Vertex (scalar toy theory)

[MQH, Kern, Alkofer, Phys.Rev.D107 (2023)]



Simpler truncation:

$$\text{gluon loop}^{-1} = \text{gluon}^{-1} - \frac{1}{2} \text{gluon loop} + \text{ghost loop}$$

The diagram shows the equation: $\text{gluon loop}^{-1} = \text{gluon}^{-1} - \frac{1}{2} \text{gluon loop} + \text{ghost loop}$. The gluon loop is represented by a red circle with two external lines. The ghost loop is represented by a green dashed circle with two external lines.

- Numerically more demanding due to calculations close to cuts.
- Branch cut

Extrapolation of $\lambda(P^2)$

Extrapolation method

- Extrapolation to time-like P^2 using Schlessinger's continued fraction method (proven superior to default Padé approximants) [Schlessinger, Phys.Rev.167 (1968)]
- Average over extrapolations using subsets of points for error estimate

$$f(x) = \frac{f(x_1)}{1 + \frac{a_1(x-x_1)}{1 + \frac{a_2(x-x_2)}{1 + \frac{a_3(x-x_3)}{\dots}}}}$$

Coefficients a_i can be determined such that $f(x)$ exact at x_i .

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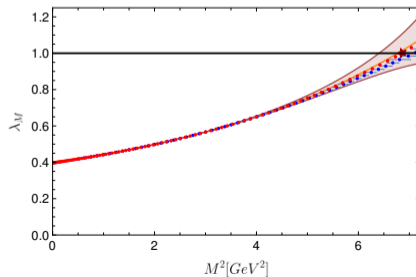
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Test extrapolation for solvable system:

Heavy meson [MQH, Sanchis-Alepuz, Fischer, Eur.Phys.J.C 80 (2020)]

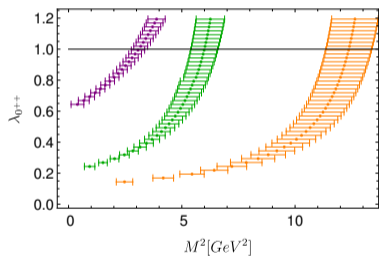
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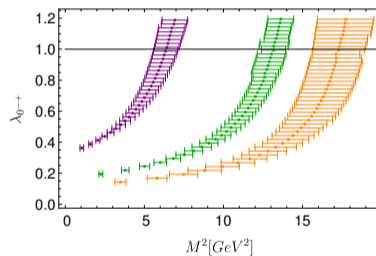


Extrapolation for glueball eigenvalue curves

0^{++} :



0^{-+} :



Several curves: ground state and excited states.

Gauge invariance

[MQH, Phys. Rev. D 101 (2020)]

Couplings can be extracted from each vertex.

- Slavnov-Taylor identities (gauge invariance): Agreement perturbatively (UV) necessary.
[Cyrol et al., Phys.Rev.D 94 (2016)]
- Difficult to realize: Small deviations \rightarrow Couplings cross and do not agree.
- Here: Vertex couplings agree down to GeV regime (IR can be different).

