The dual superconductor picture of confinement and its relation to other confinement scenarios

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June 16, 2011





Joint annual meeting of SPS and APS with SSAA and ASAA, Lausanne

Our view of the world in terms of particles

The standard model:







Our view of the world in terms of particles



MQH

Confinement

"No one has ever detected a free quark or gluon."

"Absence of free quarks and gluons."

"Quarks and gluons are not part of the physical state space."

We observe hadrons but use (unobservable) quarks and gluons as elementary fields.

Inelastic scattering:

Atoms \rightarrow constituents Hadrons \rightarrow more hadrons Without understanding confinement we do not fully understand hadron physics.

Several confinement scenarios/mechanisms

View the problem from different perspectives \rightarrow not mutually exclusive but different aspects emphasized.



see e.g. [Alkofer, Greensite, JPG34] for a short review

A few confinement scenarios and approaches

- Magnetic monopoles and dual superconductivity: ['t Hooft, Mandelstam, di Giacomo, Suzuki, ...]
- Gribov-Zwanziger scenario: role of Gribov copies? [Gribov, Zwanziger, Sorella, Dudal, ...]
- Center vortices: topological objects [Faber, Greensite, Langfeld, Olejník, ...]
- Coulomb gauge:

[Reinhardt, Watson, Zwanziger, ...]

- Landau gauge Green functions: [Alkofer, Fischer, von Smekal, Maas, Papavassiliou, Pawlowski, ...]
- AdS/CFT correspondence:

[Maldacena, Brodsky,]

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Dual superconductor picture of confinement

Conventional type-II superconductor:

- magnetic flux squeezed into vortices
- condensation of Cooper pairs



Theories where confinement was proven: compact U(1), Georgi-Glashow model, deformed $\mathcal{N}=2$ SUSY Yang-Mills

't Hooft 1976, Mandelstam 1976

Dual superconductor:

- electric flux squeezed into vortices
- condensation of magnetic monopoles



 \rightarrow In all condensation of magnetic monopoles!



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Abelian infrared dominance

Hypothesis of Abelian infrared dominance [Ezawa, Iwazaki, PRD25 (1981)]: based on dual superconductor picture





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What is the Abelian part?

$$[A_{\mu},A_{\nu}]=0?$$



Gauge actions

Action is invariant under gauge transformations. \rightarrow (Infinitely) Many equivalent configurations. Space of field configurations \mathcal{A} is not the physical space!

Gauge transformations connect physically equivalent configurations. \rightarrow Physical space is the quotient space over the gauge group G:

$$\mathcal{A}_{phys} = \mathcal{A}/G$$

Realization: Fix a gauge.



Gauge fixing

Restriction to a hypersurface $\Gamma: f[A] = 0$ (gauge fixing condition, e.g., $\partial A = 0$)



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Freedom to choose any gauge!

- Perturbation theory: Feynman gauge convenient, propagator simple $(g_{\mu\nu}/p^2)$
- Dyson-Schwinger equations: Landau gauge, simplification of calculations. Helped us learn the method.
- Deep inelastic scattering: light-cone gauge
- Dual superconductor scenario: Abelian gauges



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(All these gauges are non-perturbatively not complete: There are still equivalent configurations on Γ . \rightarrow Gribov problem Physical space \mathcal{A}_{phys} is topological not trivial!

Theorem of Singer 1978: No unique gauge fixing with a continuous gauge fixing condition.)



The maximally Abelian gauge I

Dual superconductor picture based on Abelian symmetry. What is it for SU(3)?

Approach here: Split the gauge field into diagonal and off-diagonal parts.



The maximally Abelian gauge II

- The gauge field lives in the algebra su(3): $A_{\mu} = A_{\mu}^{r} T^{r}$ (T^{r} : generators, Gellmann matrices)
- Some of them (λ_3, λ_8) are diagonal. \rightarrow Form maximal Abelian subalgebra (Cartan subalgebra). \rightarrow "Abelian"/diagonal part: A^3_{μ}, A^8_{μ}
- Consider the two parts separately

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- Consider the two parts separately \rightarrow diagonal and off-diagonal fields.
- We use the gauge freedom to choose minimized off-diagonal fields.

 $\Rightarrow \boldsymbol{D}_{\mu} \boldsymbol{A}_{\mu} = 0$

ightarrow maximally Abelian gauge (MAG)

(Compare Landau gauge: $\partial_{\mu}A_{\mu} = 0$)



The maximally Abelian gauge III

Indication of Abelian infrared dominance:

 \rightarrow String tension of quark potential stays almost the same when calculated from diagonal fields only.

[e.g. Suzuki, Yotsuyanagi, PRD42].

Note: Here the off-diagonal part is kept!

 \rightarrow Investigation of the full theory!



Functional equations

<u>Green functions:</u> Propagators and vertices \leftrightarrow describe how fields propagate and interact.

Exact relations between Green functions given by, e.g., Dyson-Schwinger and functional renormalization group equations.

DSE for the gluon propagator (Landau gauge):

$$i$$
 j -1 $+$ i j -1 $-\frac{1}{2}$ i j $-\frac{1}{2}$ j $-\frac{1}{2}$ j

Valid **non-perturbatively**: dressed propagators and vertices!

Infinite tower of coupled equations: every eq. contains higher Green functions.

Applications of functional equations

- condensed matter
- Yang-Mills theories
- QCD and QCD-like theories
- standard model physics
- supersymmetry
- gravity
- . . .
- \rightarrow Wide variety of theories.



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Lagrangian for Yang-Mills theory in the max. Abelian gauge

First example: Landau gauge





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Lagrangian for Yang-Mills theory in the max. Abelian gauge



High number of interactions \rightarrow automated derivation very useful.

 \rightarrow *Mathematica* application *DoFun* (Derivation of functional eqs.) [Alkofer, MQH, Schwenzer, CPC180; MQH, Braun, arXiv:1102.5307]

Derivation of full Dyson-Schwinger and functional RG equations.

 \rightarrow More complicated theories become accessible,

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For example: Off-diagonal gluon



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Derivation of full Dyson-Schwinger and functional RG equations.

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For example: Ghost

FSU Jena



Infrared solution for the maximally Abelian gauge I

Infinite towers of functional equations.

But: In the IR the untruncated towers can be solved qualitatively.

[MQH, Schwenzer, Alkofer, EPJC68]

Scaling-type solution: All dressing functions characterized by power laws in the IR:

 $c_i(p_1^2, p_2^2, \ldots) \stackrel{I\!R}{=} d_i \cdot (p^2)^{\kappa_i}$ (common momentum scale p)

The exponents κ_i are uniquely related!

Numerical values for the κ_i can be calculated.



Infrared solution for the maximally Abelian gauge II

Qualitative behavior of propagators at low momentum p ($\kappa_{MAG} \ge 0$):

canonical dim.

- Off-diagonal gluon propagator is IR suppressed $\sim (p^2)^{\kappa_{MAG}-1}$.
- Ghost propagator (off-diag.) is IR suppressed $\sim (p^2)^{\kappa_{MAG}-1}$.
- Diagonal gluon propagator is IR enhanced $\sim (p^2)^{-\kappa_{MAG}-1}$. \rightarrow Diagrams with most diagonal gluons dominate in DSEs.
- Qualitative solution for the whole tower of Green functions.

 \rightarrow IR dominance of diagonal/"Abelian" degrees of freedom.



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(SU(N > 2) has more interactions than SU(2), but the IR solution is the same.)



Value of the infrared exponent κ_{MAG}

Solution for κ_{MAG} is <u>necessary but not sufficient</u>. Dressing functions of gluons and ghosts:

$$\begin{split} c_{A}(p^{2}) \stackrel{p^{2} \to 0}{=} d_{A} \cdot (p^{2})^{-\kappa_{MAG}} & c_{c}(p^{2}) \stackrel{p^{2} \to 0}{=} d_{c} \cdot (p^{2})^{\kappa_{MAG}} \\ c_{B}(p^{2}) \stackrel{p^{2} \to 0}{=} d_{B} \cdot (p^{2})^{\kappa_{MAG}} & 0 \leq \kappa_{MAG} \leq 1 \end{split}$$

 $\kappa_{\textit{MAG}}\approx 0.74$

 \rightarrow Infrared consistent solution exists.

Gauge fixing condition: $D_{\mu}A_{\mu} = 0 \leftrightarrow$ gauge fixing parameter α_{MAG} , MAG: $\alpha_{MAG} = 0$ (cf. linear covariant gauges: $\partial_{\mu}A_{\mu} = 0 \leftrightarrow \alpha_{LCG}$, Landau gauge: $\alpha_{LCG} = 0$)



Gauge fixing parameter dependence

Extension to non-zero α_{MAG} easily possible, not easy for Landau gauge and linear covariant gauges.

Two solution branches:



General linear covariant gauges: different from the Landau gauge \rightarrow The maximally Abelian is the first case where an IR solution is found that seems independent of the gauge fixing parameter.



The maximally Abelian and the Landau gauge

Structures of equations considerably different in the two gauges.



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Understanding of confinement increases: We put together the jigsaw pieces, but the big picture needs to be finished.

- Maximally Abelian gauge is the second gauge after the Landau gauge where a qualitative solution for the whole tower of Green functions is known.
- Relation to Landau gauge partly understood:

 $\mathsf{ghost} \leftrightarrow \mathsf{Abelian} \ \mathsf{dominance}$

- Maximally Abelian gauge suited for investigation of Abelian IR dominance.
- Results corroborate this scenario, applicable for family of gauges (α_{MAG}).



Acknowledgments

Thanks to the following people

for their collaboration, support and helpful discussions:

Reinhard Alkofer

Kai Schwenzer, Silvio P. Sorella Christian S. Fischer, Axel Maas, Jan M. Pawlowski, Lorenz von Smekal, Daniel Zwanziger David Dudal, Leonard Fister, Marcelo S. Guimarães, Klaus Lichtenegger, Nele Vandersickel

This work was supported by







