Going beyond the propagators of Landau gauge Yang-Mills theory

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 $X_{th}\ Confinement$ and the Hadron Spectrum

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Non-perturbative Landau gauge Green functions

Non-perturbative propagators of Landau gauge Yang-Mills theory:

- information about confinement
- input for phenomenological calculations (QCD phase diagram, bound states, ...) → sessions B, D
- QCD phase diagram:

phase transitions (e.g. via condensates, Polyakov loop), no sign problem for functional methods, but equations get more complicated;

also vertices change, e.g., ghost-gluon vertex at non-zero temperature [Fister, Pawlowski, 1112.5440]

Calculated with

- Monte-Carlo simulations,
- functional renormalization group,
- Dyson-Schwinger equations,

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Dyson-Schwinger equations: Propagators

Dyson-Schwinger equations (DSEs) of gluon and ghost propagators:



- Equations of motion of correlation functions.
- Infinite tower of coupled integral equations.
- Contain three-point and four-point functions: ghost-gluon vertex, three-gluon vertex, four-gluon vertex

Dyson-Schwinger equations: Propagators

Dyson-Schwinger equations (DSEs) of gluon and ghost propagators:



Truncated propagator Dyson-Schwinger equations

Standard truncation:



using bare ghost-gluon vertex and three-gluon vertex model

Truncated propagator Dyson-Schwinger equations

Standard truncation:





- Various lattice results [Cucchieri, Maas, Mendes, PRD77; Ilgenfritz, BJP37]
- OPE analysis [Boucaud et al., JHEP 1112]
- Modeling via ghost DSE [Dudal, Oliveira, Rodriguez-Quintero, 1207.5118]
- Semi-perturbative DSE analysis [Schleifenbaum et al., PRD72]

$$\Gamma^{A\bar{c}c,abc}_{\mu}(k;p,q) := ig f^{abc} \left(p_{\mu} A(k;p,q) + k_{\mu} B(k;p,q) \right)$$

Note:

B(k; p, q) is irrelevant in Landau gauge, but it is not the pure longitudinal part.

IR and UV consistent truncation:



Solutions of DSEs: Decoupling and scaling

- Two types of solutions with functional methods that differ only in deep IR [Boucaud et al., JHEP 0806, 012; Fischer, Maas, Pawlowski, AP 324]: scaling [von Smekal, Alkofer, Hauck PRL97], decoupling [Aguilar, Binosi, Papavassiliou PRD78]
- Lattice calculations find only decoupling type solution for d = 3, 4and scaling for d = 2



d = 2: Analytic and numerical arguments from DSEs for scaling only [Cucchieri, Dudal, Vandersickel, PRD85; MQH, Maas, von Smekal, 1207.0222] as well as from analysis of Gribov region [Zwanziger, 1209.1974].

Ghost-loop-only truncation

Ghost-loop-only truncation

Scaling solution: Ghost loop is IR dominant by power counting [von Smekal, Hauck, Alkofer, 79; Alkofer, Fischer, Llanes-Estrada, PLB611] Decoupling solution: No ghosts \rightarrow no gluon "mass" [Aguilar, Binosi, Papavassiliou, JHEP1201]; ghost loop alone, i.e., $\Gamma^{A^3} = 0$? $G(p^2)$ $Z(p^2)/p^2$ 3.0 2.5 Both loops 1.5 Both loops 2.0Ghost loop only Ghost loop only 1.0 1.5 0.5 1.0 p^{2} $p_{10^4}p$ 0.5 0.0 10-4 0.01 100 10^{-4} 0.01 100

 \Rightarrow Ghost loop alone sufficient to obtain decoupling behavior!

Note: Gluon loop only \rightarrow Mandelstam solution, IR divergent gluon propagator [Mandelstam PRD20]

Three-gluon vertex: Ultraviolet

Ansatz that reproduces the correct UV behavior of the gluon propagator [Fischer, Alkofer, Reinhardt, PRD65]:

$$D^{A^{3}}(x, y, z) = \frac{1}{Z_{1}} \frac{[G(y)G(z)]^{1-a/\delta-2a}}{[Z(y)Z(z)]^{1+a}}$$

y and z are the momenta in the gluon loop, i.e., not Bose symmetric in y, z and x.

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Bose symmetrized version:

$$D^{A^3,UV}(x,y,z) = \frac{1}{Z_1}G(x+y+z)^{\alpha}Z(x+y+z)^{\beta}$$

Fix α and β :

- 1 UV behavior of gluon propagator
- 2 IR behavior of three-gluon vertex?

Three-gluon vertex: Ultraviolet

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Fix α and β :

- 1 UV behavior of gluon propagator
- 2~ IR behavior of three-gluon vertex \rightarrow yes, but \ldots

Three-gluon vertex: Infrared

Hints from lattice data [Cucchieri, Maas, Mendes, PRD77]:

Three-gluon vertex might have a zero crossing.

(d=2,3: zero crossing seen [Cucchieri, Maas, Mendes, PRD77; Maas, PRD75])

Three-gluon vertex: Infrared

Hints from lattice data [Cucchieri, Maas, Mendes, PRD77]:

Three-gluon vertex might have a zero crossing. (d = 2, 3: zero crossing seen [Cucchieri, Maas, Mendes, PRD77; Maas, PRD75])

 \Rightarrow add IR part:

$$D^{A^{3},IR}(x,y,z) = h_{IR}G(x+y+z)^{3}(f^{3g}(x)f^{3g}(y)f^{3g}(z))^{4}$$

IR damping function $f^{3g}(x) := \frac{\Lambda_{3g}^2}{\Lambda_{3g}^2 + x}$

Parameter for zero crossing: $h_{IR} < 0$

New three-gluon vertex:

$$D^{A^{\mathbf{3}}}(x,y,z)=D^{A^{\mathbf{3}},IR}(x,y,z)+D^{A^{\mathbf{3}},UV}(x,y,z)$$

Propagators: Improving the three-gluon vertex

ghost-gluon vertex: bare



original three-gluon vertex

Propagators: Improving the three-gluon vertex

ghost-gluon vertex: bare



original three-gluon vertex Bose symmetric three-gluon vertex

Propagators: Improving the three-gluon vertex

ghost-gluon vertex: bare



original three-gluon vertex Bose symmetric three-gluon vertex Bose symmetric three-gluon vertex with IR part

 \Rightarrow Improved three-gluon vertex adds additional strength in the mid-momentum regime.

$$\Gamma^{Aar{c}c,abc}_{\mu}(k;p,q) \coloneqq i\,g\,f^{abc}\left(p_{\mu}A(k;p,q)+k_{\mu}B(k;p,q)
ight)$$

basis choice: p^2 , k^2 , ϕ

$$\Gamma^{A\bar{c}c,abc}_{\mu}(k;p,q) := i g f^{abc} \left(p_{\mu} A(k;p,q) + k_{\mu} B(k;p,q) \right)$$



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MQH

Ghost-gluon vertex: Selected configurations

Fixed angle:



Fixed momentum:



Results for SU(2)

Ghost-gluon vertex: Comparison with lattice data

Symmetric configuration $k^2 = q^2 = p^2$:



Orthogonal configuration $k^2 = 0$, $q^2 = p^2$:



lattice data: [Cucchieri, Maas, Mendes, PRD77]

Results for SU(2)

Propagators and dynamic ghost-gluon vertex



Propagators and dynamic ghost-gluon vertex





bare ghost-gluon vertex, Bose non-symmetric three-gluon vertex bare ghost-gluon vertex, improved three-gluon vertex, improved three-gluon vertex lattice: 32⁴ lattice: 48⁴ [Sternbeck, hep-lat/0609016]

Propagators and dynamic ghost-gluon vertex

- More realistic three-gluon vertex improves mid-momentum behavior.
- Dynamic ghost-gluon vertex influences ghost in IR Note: Solution with dynamic ghost-gluon vertex requires improved three-gluon vertex for stability (gluon loop otherwise too strong).
- Missing part around 2 GeV probably due to two-loop terms
 talk of Valentin Ma

 \rightarrow talk of Valentin Mader



Summary & Conclusions

- Ghost dominance for decoupling and scaling solution (ghost-loop-only analysis).
- More realistic three-gluon vertex: Bose symmetric and IR features from lattice
 → Yields more support in mid-momentum regime. → Reduces gap to lattice results.
- Dynamic inclusion of ghost-gluon vertex:
 - rather small changes for gluon
 - influences ghost in the IR
 - improved three-gluon vertex required
- Ghost-gluon vertex results agree well with available lattice results.

Summary & Conclusions

• Inclusion of three-point functions:

- test of truncations,
- required for quantitative results and
- likely also for some aspects of non-zero temperature and density calculations
- Automatization tools:

DoFun [Alkofer, MQH, Schwenzer, CPC180; MQH, Braun, CPC183] CrasyDSE [MQH, Mitter, CPC183]

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Thank you very much for your attention.