

Quantum chromodynamics from the functional point of view



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Outline

- ① Introduction
- ② Functional equations
- ③ Dyson-Schwinger equations
- ④ DSEs in QCD
- ⑤ Results in the Yang-Mills sector

The world in terms of particles

The standard model:

	d	s	b	g
Quarks	u	c	t	γ
Leptons	e	μ	τ	Z
	ν_e	ν_μ	ν_τ	W^\pm
Bosons				

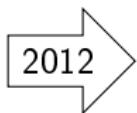
Higgs?

The world in terms of particles

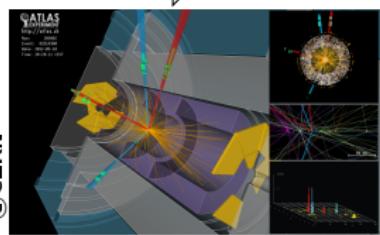
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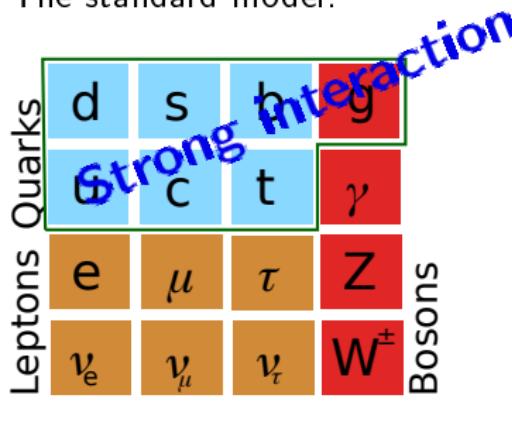


Which Higgs?



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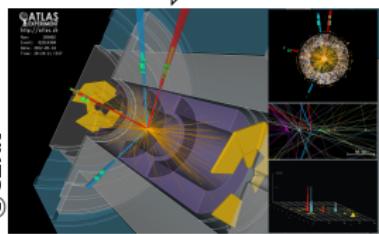
The standard model:



Higgs?

2012

Which Higgs?



Quantum chromodynamics:
self-consistent by itself, could be even
fundamental

The world in terms of particles

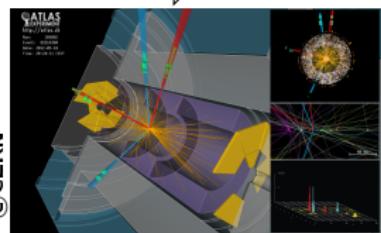
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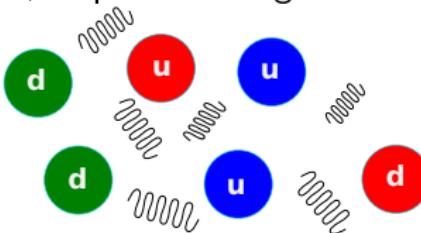
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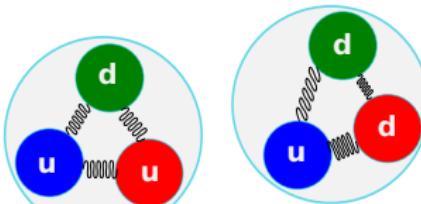


fundamental fields: quarks and gluons

u, d quarks are light \sim MeV



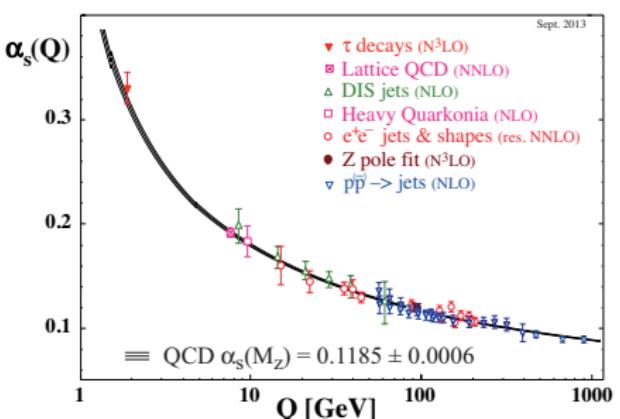
physical degrees of freedom: hadrons \sim GeV



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How to investigate QCD

Asymptotic freedom (Nobel prize 2004):

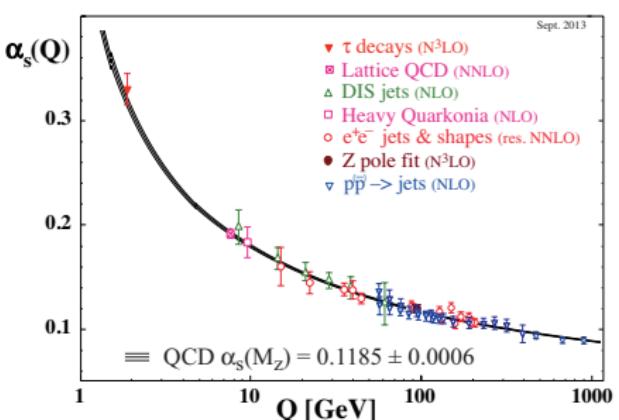


Perturbative description at high energies.

Plenty of applications.

How to investigate QCD

Asymptotic freedom (Nobel prize 2004):



Perturbative description at high energies.
Plenty of applications.

- Perturbative series is not convergent.
 - Non-perturbative phenomena?
 - E.g., no mass creation to every order in perturbation theory.

⇒ Non-perturbative methods required.

Perturbation theory based on non-perturbative 'models', e.g.,

- (Refined) Gribov-Zwanziger model
 - Massive extension [Peláez, Reinosa, Serreau, Tissier, Tresmontant, Wschebor, '10-'14]

The family of functional equations

Coupled integro-differential/integral equations.

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Coupled integro-differential/integral equations.

- Dyson-Schwinger equations: eqs. of motion for correlation functions

$$\begin{array}{c} \text{---} \xrightarrow[S(p)]{-1} \text{---} \\ \text{---} \xrightarrow[S_0(p)]{-1} \text{---} + \end{array} \quad \begin{array}{c} \text{---} \xrightarrow[\gamma_\mu \text{---}]{} \text{---} \xrightarrow[S(q)]{\Gamma_\mu(p,q)} \text{---} \end{array}$$

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$$-\overrightarrow{S(p)} \xrightarrow{-1} = -\overrightarrow{S_0(p)} \xrightarrow{-1} + \gamma_\mu \text{ (loop diagram)} \xrightarrow{\Gamma_\mu(p,q)} S(q)$$

- Functional renormalization group: flow equations, RG scale k , regulator

$$k \frac{\partial}{\partial k} \text{---} \bullet^{-1} = - \text{---} \circlearrowleft + \frac{1}{2} \text{---} \circlearrowright$$

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- N-PI effective action

The family of functional equations

Coupled integro-differential/integral equations.

- Dyson-Schwinger equations: eqs. of motion for correlation functions

$$-\overrightarrow{S(p)} \xrightarrow{-1} = -\overrightarrow{S_0(n)} \xrightarrow{-1} + \gamma_\mu \begin{array}{c} \text{Diagram of a loop with } S(a) \text{ and } \Gamma_\mu(p, q) \\ \text{with a dot at the top vertex} \end{array}$$

- Functional renormalization group: flow equations, RG scale k , regulator

$$k \frac{\partial}{\partial k} \text{---} \bullet^{-1} = - \text{---} \circlearrowleft + \frac{1}{2} \text{---} \circlearrowright$$

- N-PI effective action

Non-perturbative in the sense:

- Exact equations.
 - No small coupling required.

In reality they cannot be solved exactly (with a few exceptions).
Self-consistence!

Comparison: DSEs and flow equations

Dyson-Schwinger equations (DSEs)	Functional RG equations (FRGEs)
'integrated flow equations'	'differential DSEs'
effective action $\Gamma[\phi]$	effective average action $\Gamma^k[\phi]$
-	regulator
n-loop structure ($n \text{ const.}$)	1-loop structure
exactly only bare vertex per diagram	no bare vertices
$\frac{\partial}{\partial \phi} \Gamma[\phi] = \text{Diagrammatic expansion}$	$k \frac{\partial}{\partial k} \Gamma^k[\phi] = \text{Diagrammatic expansion}$

- Both **systems of equations** are **exact**.
 - Both contain infinitely many equations.

From Green functions to 'observables'

Functional equations are expressed in terms of Green functions/correlation functions/n-point functions.

The effective action is the generating functional of 1PI Green functions.

2

The set of **all** Green functions describes the theory completely.

Green functions \Rightarrow 'observables'?

Examples:

- Bound state equations → masses and properties of hadrons
 - Analytic properties of Green functions → confinement
 - (Pseudo-)Order parameters

Functional equations and lattice methods

	functional equations	lattice
source of error	truncation	finite volume, finite lattice spacing, statistics
temperature	✓	✓
chemical potential	✓	sign problem
analytic structure	✓	no

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Green functions	functional equations	lattice
propagators	✓	✓
three-point functions	ghost-gluon vertex: ✓ 3-gluon vertex: ✓ quark-gluon vertex: (✓)	limited mom. dependence
four-point functions	✓	not soon

Dyson-Schwinger/Schwinger-Dyson-equations

- ① Start from path integral: Integral of derivative vanishes.

$$0 = \int D[\phi] \frac{\delta}{\delta \phi} e^{-S + \int dy \phi(y) J(y)}$$

Details and example of scalar theory:

<http://tinyurl.com/dsenotes>

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 - ③ Master equation:

$$\frac{\delta S}{\delta \phi(x)} \Bigg|_{\phi(x') = \phi_{\text{cl}}(x') + \int dz D(x', z) J} = \frac{\delta \Gamma[\phi_{\text{cl}}]}{\delta \phi_{\text{cl}}(x)}$$

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- ④ DSEs for Green functions by differentiating wrt fields.

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Automated derivation

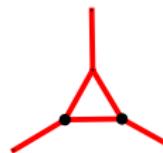
Derivation by hand becomes tedious:

- Large Lagrangians.
 - Higher Green functions.
 - Larger truncations.
 - Error-prone.

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$$\begin{aligned} & \left(2 q^2 Nc Z1 DAAA \left[y, qs + y + 2 sp[q, q1], \frac{-y - sp[q, q1]}{\sqrt{y (qs + y + 2 sp[q, q1])}} \right] \right. \\ & \left. DAAA \left[x2 + y + 2 sp[p, q], qs + x2 - 2 sp[p, q1], \frac{-x2 - sp[p, q] + sp[p, q1] + sp[q, q1]}{\sqrt{(x2 + y + 2 sp[p, q]) (qs + x2 - 2 sp[p, q1])}} \right] Dg1[qs] Dg1[qs + x2 - 2 sp[p, q1]] Dg1[\right. \\ & sp[p, q]^4 (sp[p, q1]^2 sp[q, q1] (y + sp[q, q1]) + qs x2 (y (9 qs + 6 (x2 + y)) + (5 qs + 6 x2 + 10 y) sp[q, q1]) - sp[p, q1] (qs y (5 qs + sp[p, q]^3 (2 sp[p, q1]^3 (qs y - sp[q, q1]^2) + sp[p, q1] (qs y (10 qs^2 + (-5 x2 - 3 y) y + qs (19 x2 + 3 y)) + (3 qs^3 + 8 qs x2 y + 21 qs^2 + qs x2 (y (-9 qs^2 + 3 x2^2 + 7 x2 y + 3 y^2 + 2 qs (x2 + y)) + (-10 qs^2 + qs (-3 x2 - 19 y) x2 (3 x2 + 5 y)) sp[q, q1] + (-16 qs - 7 x2 - 11 y) sp[p, q1]^2 (qs (-16 qs - 11 x2 - 7 y) y + (-5 qs^2 + qs (-9 x2 - 19 y) + 2 y (5 x2 + 3 y)) sp[q, q1] + (-5 qs + 12 (x2 + y)) sp[q, q1]^2 + sp[p, q]^2 (sp[p, q1]^4 sp[q, q1] (qs + sp[q, q1]) + sp[p, q1]^3 (qs y (7 qs + 11 x2 + 16 y) + (-6 qs^2 + y (9 x2 + 5 y) + qs (-10 x2 + 19 y) qs x2 (y (-3 qs^3 - 10 qs^2 (x2 + y) - 6 x2 y (x2 + y) + qs (-3 x2^2 - 19 x2 y - 3 y^2)) + (-6 qs^3 + qs^2 (-21 x2 - 32 y) + qs (-9 x2^2 - 60 x2 y (-15 qs^2 - 15 x2^2 + qs (-46 x2 - 41 y) - 41 x2 y - 12 y^2) sp[q, q1]^2 + (-7 qs - 16 x2 - 11 y) sp[q, q1]^3) + sp[p, q1]^2 (qs y (-15 qs^3 + qs^2 (5 x2 - 39 y) + qs (-81 x2 - 39 y) y + y^2 (5 x2 + 3 y)) sp[q, q1] + (12 qs^2 + 12 x2^2 + 3 x2 y + 12 y^2 + 3 qs (x2 + y)) sp[q, q1] (qs y (6 qs^3 + qs^2 (32 x2 + 21 y) + qs (25 x2^2 + 60 x2 y + 9 y^2) + x2 (3 x2^2 + 25 x2 y + 15 y^2)) + (15 qs^3 (x2 + y) + x2 y (-3 x2^2 + x2 y (-3 qs^3 + x2^2 (-3 x2 - 5 y) + qs^2 (39 x2 - 5 y) + qs x2 (39 x2 + 81 y)) sp[q, q1]^2 + (-6 qs^2 + qs (19 x2 - 10 y) + x2 (5 x2 + 9 y)) sp[q, q1] (x2 y (-sp[p, q1]^5 (qs + sp[q, q1]) + sp[p, q1]^4 (qs (6 qs + 6 x2 + 9 y) + (10 qs + 6 x2 + 5 y) sp[q, q1]) - qs (qs y - sp[q, q1]^2) (x2 (-6 qs + 9 x2 + 6 y) sp[q, q1]^2 + sp[q, q1]^3) + sp[p, q1]^3 (qs (-3 qs^2 - 3 x2^2 + qs (-7 x2 - 2 y) - 2 x2 y + 9 y^2) + (-3 x2^2 + 3 x2 y + sp[p, q1]^2 (qs (-3 qs^2 (2 x2 + y) + qs (-6 x2^2 - 19 x2 y - 10 y^2) + y (-3 x2^2 - 10 x2 y - 3 y^2)) + (-3 qs^3 - 25 qs^2 (x2 + y) + qs (-15 x2^2 - 15 x2^2 - 46 x2 v - 15 v^2 - 41 qs (x2 + v)) sp[q, q1] (alpha_11^2 + (-11 alpha_16 x2 - 7 v) sp[q, q1]^3) + \right. \end{aligned}$$

Automated derivation

Derivation by hand becomes tedious:

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- Automated derivation:

Mathematica package *DoFun* [Alkofer, Huber, Schwenzer '08; Huber, Braun '11]

<http://tinyurl.com/dofun2>

- Framework for numeric handling:

C++ program *CrasyDSE* [Huber, Mitter '11]

<http://tinyurl.com/crasydse>

Is it convergent?

Often people tend to think 'perturbatively':

No small parameter. → What means to control the calculation?

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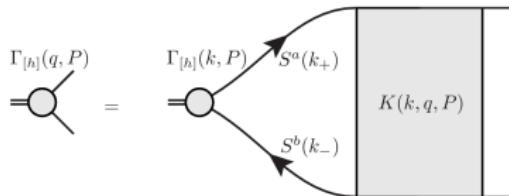
No small parameter. \rightarrow What means to control the calculation?

- Comparisons with
 - perturbation theory
 - lattice calculation
 - Sometimes analytic results possible.
 - Deform truncation.

Will come back to this for Yang-Mills theory.

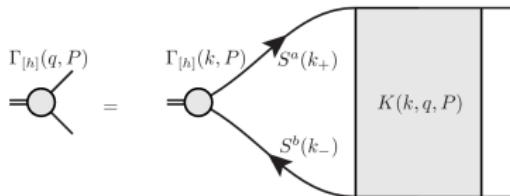
Bound state equations (BSE)

Example: Mesons from Bethe-Salpeter equation (BSE) with rainbow-ladder approximation

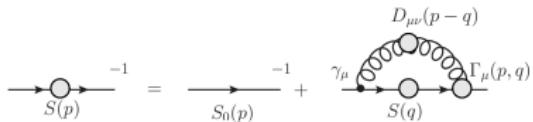


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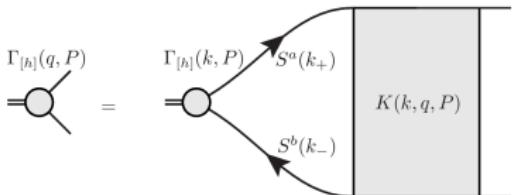


Contains quark propagator S and kernel K .

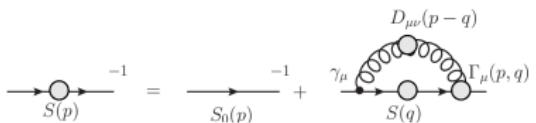


Bound state equations (BSE)

Example: Mesons from Bethe-Salpeter equation (BSE) with rainbow-ladder approximation



Contains quark propagator S and kernel K .



Rainbow-ladder

- $K \rightarrow$ dressed one gluon exchange
 - effective gluon propagator
 - bare quark-gluon vertex

Rainbow-ladder approximation

What is special about rainbow-ladder?

Respects axial-vector Ward-Takahashi identity.

→ Consistent with chiral symmetry.

→ Pion is the massless Goldstone boson in the chiral limit.

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Dressed one-gluon exchange parametrized by model for the coupling:

Maris-Tandy '99, one effective parameter

$$\frac{G(p^2)}{p^2} = \frac{4\pi^2 D}{\omega^6} p^2 \exp^{-p^2/\omega^2} + 4\pi \frac{\gamma_m \pi F(p^2)}{1/2 \ln(\tau + (1 + p^2/\Lambda_{\text{QCD}})^2)}$$

- Good description of many mesons with light quarks.

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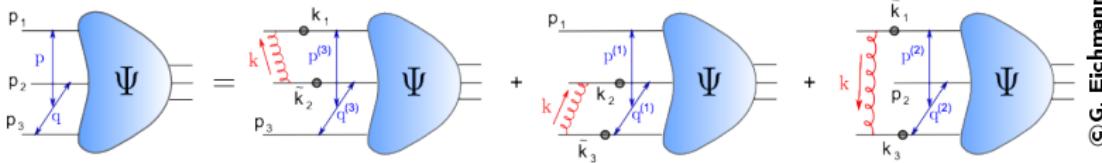
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Baryons: Faddeev equation

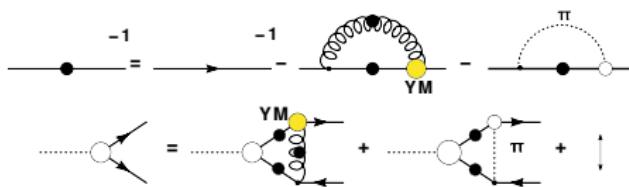


Beyond rainbow-ladder: Examples

- Vertex dressing [Bender, Roberts, von Smekal '96]:



- Include pion back coupling effects [e.g., Fischer, Nickel, Wambach '07; Fischer, Nickel, Williams '08; Fischer, Williams '08]:



- Include gluon self-interaction [e.g., Maris, Tandy '06; Fischer, Williams '09]

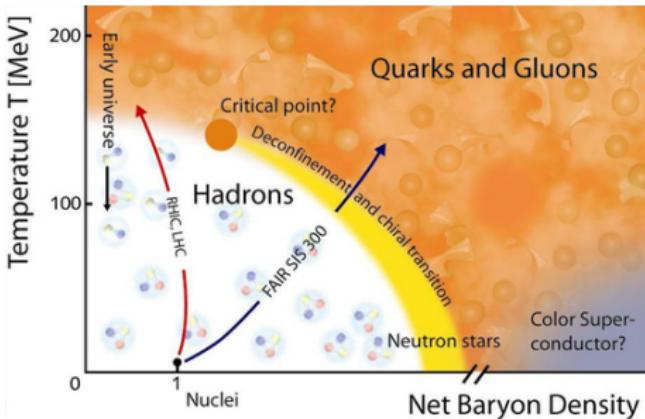


→ Three-gluon vertex required!

QCD phase diagram

Examples for accessible quantities:

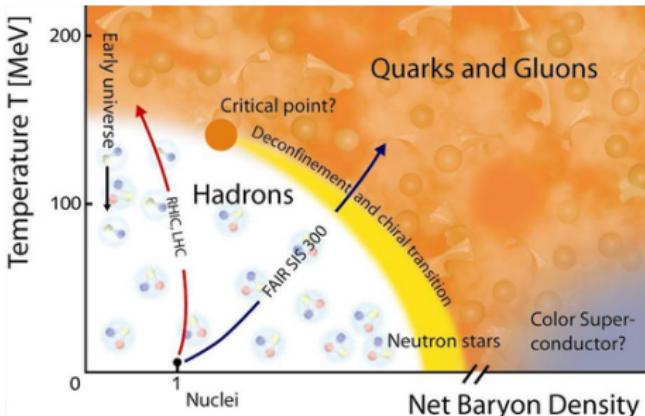
- quark condensate (chiral symm.)
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QCD phase diagram

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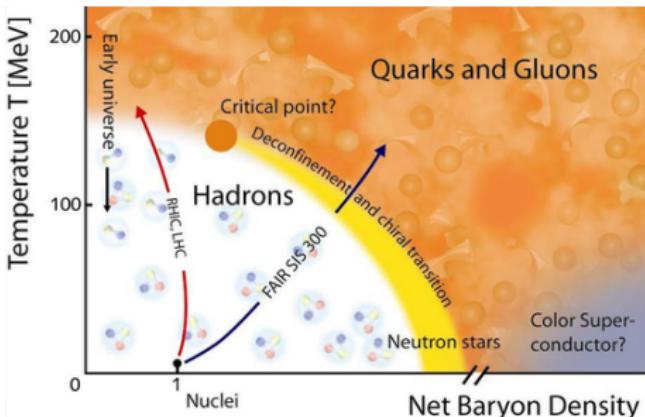
Input

- Functional equations: propagators, three-point functions

QCD phase diagram

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Input

- Functional equations: propagators, three-point functions
 - Lattice: propagators at $\mu = 0$, three-point functions
 - Models

Landau gauge Yang-Mills theory

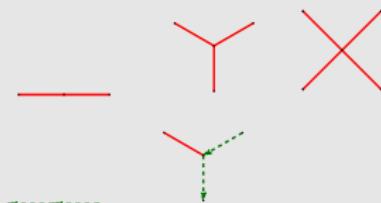
Gluonic sector of quantum chromodynamics: Yang-Mills theory

$$\mathcal{L} = \frac{1}{2} F^2 + \mathcal{L}_{gf} + \mathcal{L}_{gh}$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu + i g [A_\mu, A_\nu]$$

Landau gauge

- simplest one for functional equations
 - $\partial_\mu \mathbf{A}_\mu = 0: \quad \mathcal{L}_{gf} = \frac{1}{2\xi} (\partial_\mu \mathbf{A}_\mu)^2, \quad \xi \rightarrow 0$
 - requires **ghost** fields: $\mathcal{L}_{gh} = \bar{\mathbf{c}} (-\square + g \mathbf{A} \times) \mathbf{c}$



The tower of DSEs

$$\begin{aligned}
 & \text{Diagram 1:} \quad i \text{---} j^{-1} = + \text{Diagram 2: } i \text{---} j^{-1} - \frac{1}{2} \text{Diagram 3: } i \text{---} j^{-1} - \frac{1}{2} \text{Diagram 4: } i \text{---} j^{-1} + \text{Diagram 5: } i \text{---} j^{-1} \\
 & \text{Diagram 2: } -\frac{1}{6} \text{Diagram 6: } j \text{---} i^{-1} - \frac{1}{2} \text{Diagram 7: } j \text{---} i^{-1} \\
 & \text{Diagram 3: } \text{gluon propagator} \\
 & \text{Diagram 4: } \text{ghost propagator}
 \end{aligned}$$

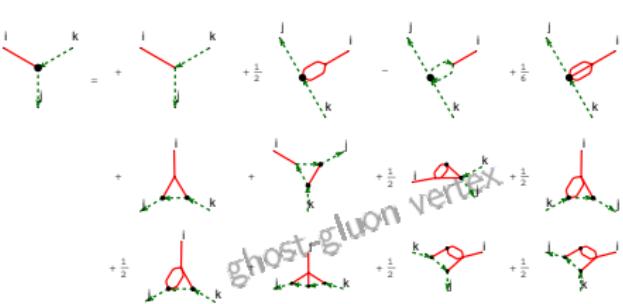
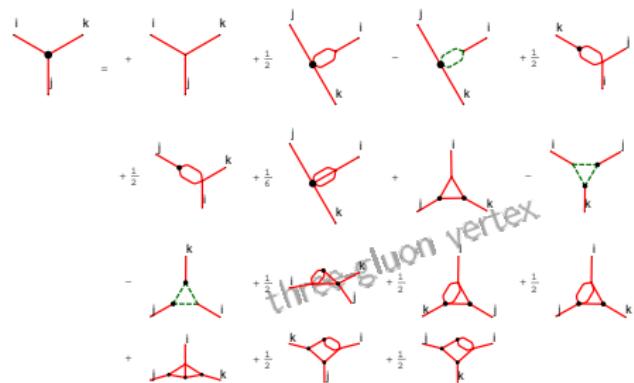
The tower of DSEs

$$i \text{---} \bullet \text{---} j = -\frac{1}{6} i \text{---} \bullet \text{---} j + \frac{1}{2} i \text{---} \circ \text{---} j - \frac{1}{2} i \text{---} \bullet \text{---} j + \frac{1}{2} i \text{---} \bullet \text{---} j + i \text{---} \bullet \text{---} j$$

gluon propagator

$$j \quad \bullet \quad i^{-1} + j \quad i^{-1} - j \quad \text{ghost propagator}$$

Self-consistency!



Infinitely many equations. In QCD, every n -point function depends on $(n + 1)$ - and possibly $(n + 2)$ -point functions.

Taming the equations

Keep most important parts!

- Drop quantities
 - Model quantities

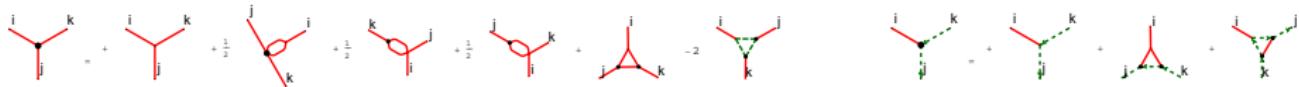
Taming the equations

Keep most important parts! The art ...

- Drop quantities
 - Model quantities

$$\text{Diagram 1} = \text{Diagram 2} - \frac{1}{2} \text{Diagram 3} + \text{Diagram 4}$$

$$j \cdot \bullet \cdot j^{-1} = j \cdot j^{-1} - j \cdot \text{red loop} \cdot j$$

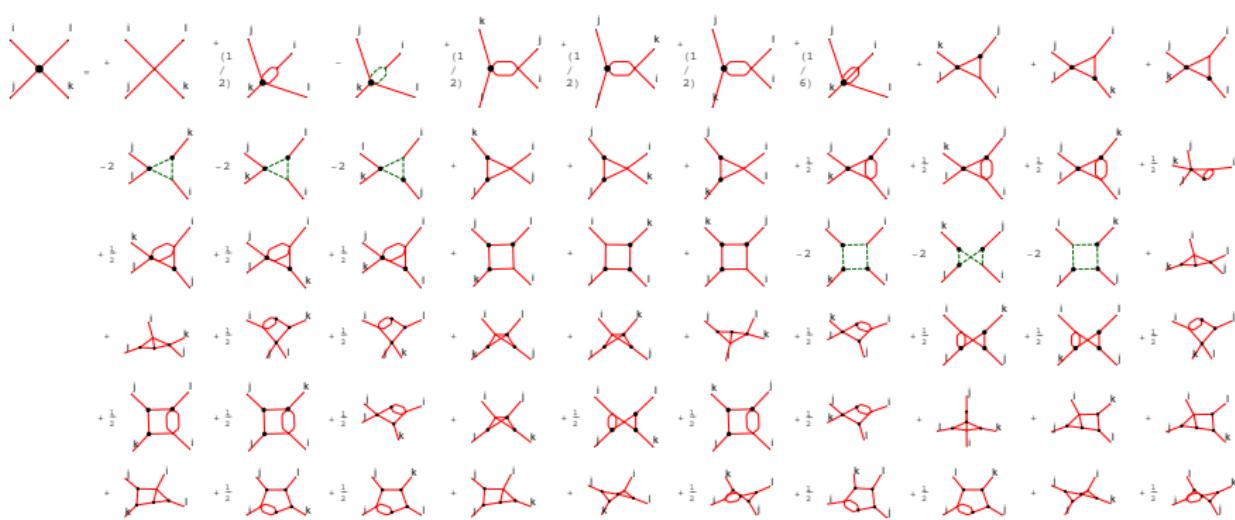


Most important parts

- UV leading (perturbation theory)
 - IR leading (analytic, lattice)

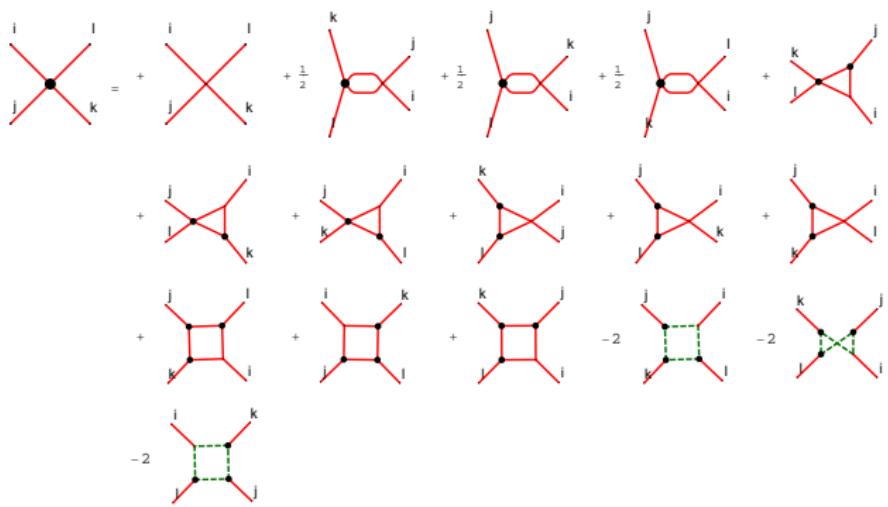
Four-gluon vertex

- 20 one-loop, 39 two-loop diagrams



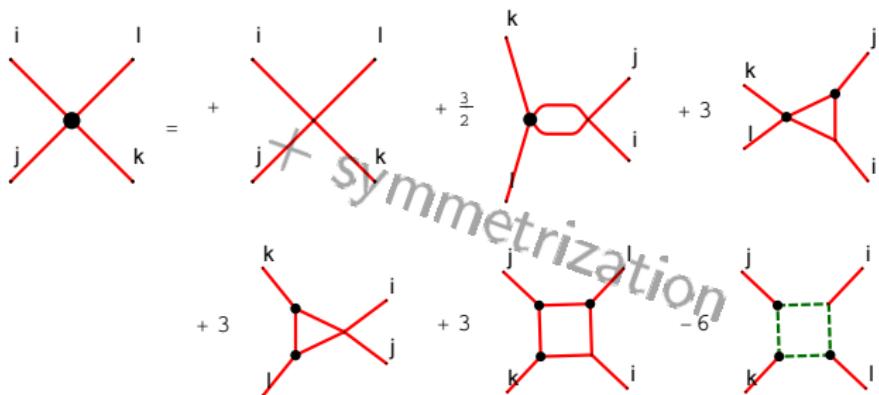
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 - Keep UV leading diagrams $\rightarrow 16$ diagrams



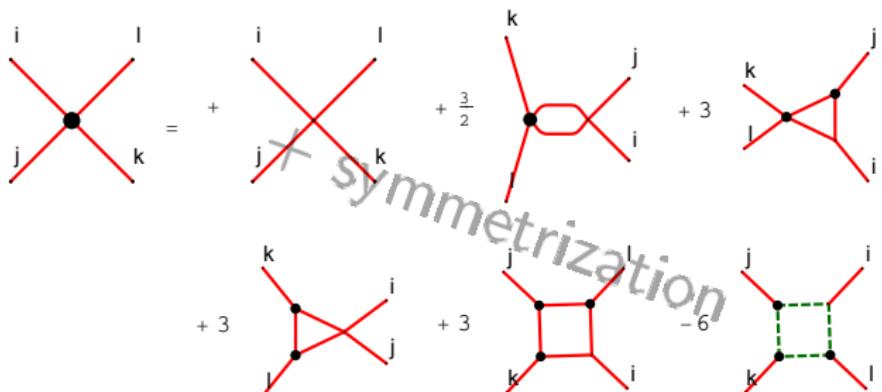
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- Calculate **full** momentum dependence.
 \rightarrow Access to all permutations of this diagram. $\rightarrow 6$ diagrams



Four-gluon vertex

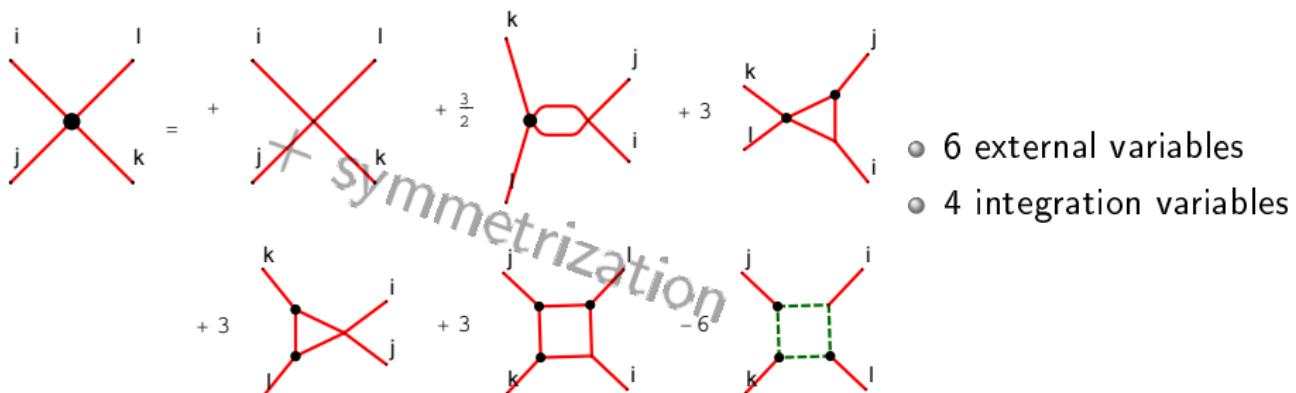
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No model dependence! \rightarrow 'Truncation closes.'

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Truncation recap

Principle

- Drop or model unknown quantities ...
- ...but capture the important part.

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Yang-Mills theory:

- Complete at leading UV order:
primitively divergent quantities only. → 5 coupled DSEs
- Truncation closes.

In the following consider subparts of this system.

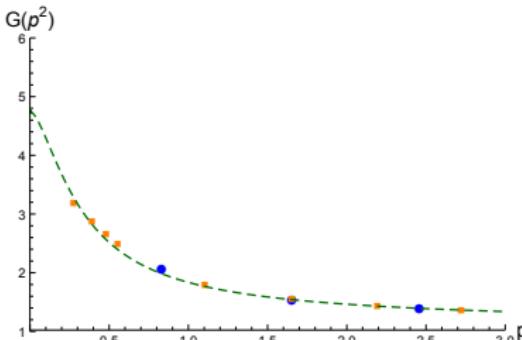
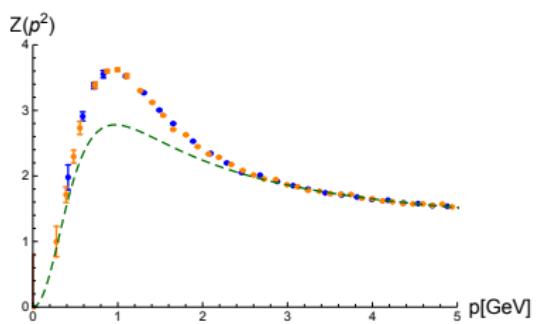
Propagators I

$$\text{Diagram: } i \text{ (red line)} - j \text{ (black dot)}^{-1} = i \text{ (red line)} - j \text{ (black dot)}^{-1} - \frac{1}{2} i \text{ (red line)} - j \text{ (green dashed line)}^{-1} + i \text{ (red line)} - j \text{ (green dashed line)}^{-1}$$

$$\text{Diagram: } j \text{ (green dashed line)} - i \text{ (black dot)}^{-1} = j \text{ (green dashed line)} - i \text{ (black dot)}^{-1} - j \text{ (green dashed line)} - i \text{ (red line)}^{-1}$$

Long-time standard truncation

- No four-gluon vertex
- Ghost-gluon vertex: bare
- Three-gluon vertex: model



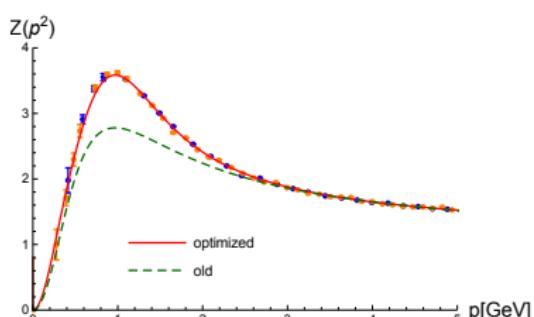
Propagators I

$$\text{---} \bullet \text{---} j^{-1} = \text{---} i \text{---} j^{-1} - \frac{1}{2} \text{---} i \text{---} \text{---} i + \text{---} i \text{---} \text{---} i$$

$$\text{---} j^{-1} = \text{---} i \text{---} j^{-1} - \text{---} j \text{---} i$$

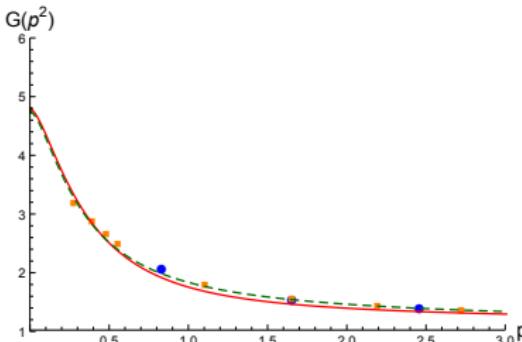
Long-time standard truncation

- No four-gluon vertex
- Ghost-gluon vertex: bare \rightarrow dressed (dynamic)
- Three-gluon vertex: model \rightarrow optimized model



[MQH, von Smekal '12; lattice; Sternbeck '06]

→ Role of three-gluon vertex?



→ Use as input in other calculations.

Propagators I

$$\text{---} \bullet \text{---} = \text{---} \text{---} \text{---}^{-1} - \frac{1}{2} \text{---} \text{---} \text{---}$$

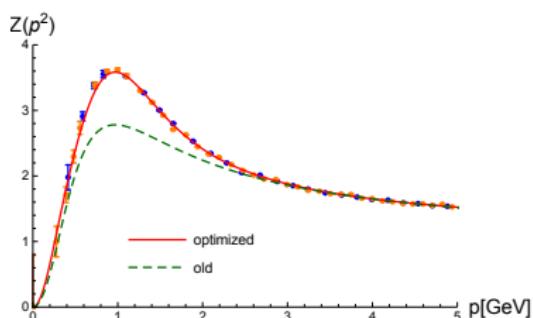
$$\text{---} \bullet \text{---}^{-1} = \text{---} \text{---} \text{---}^{-1} - \text{---} \text{---} \text{---}$$

Long-time standard truncation

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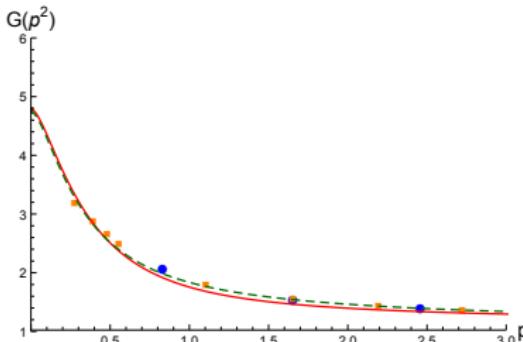
Missing strength in mid-momentum regime:

- neglected diagrams?
- vertices?



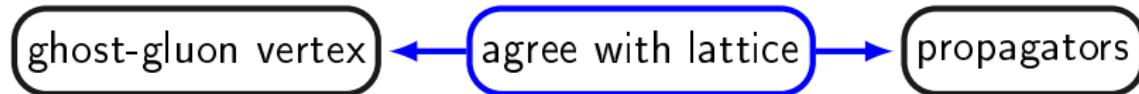
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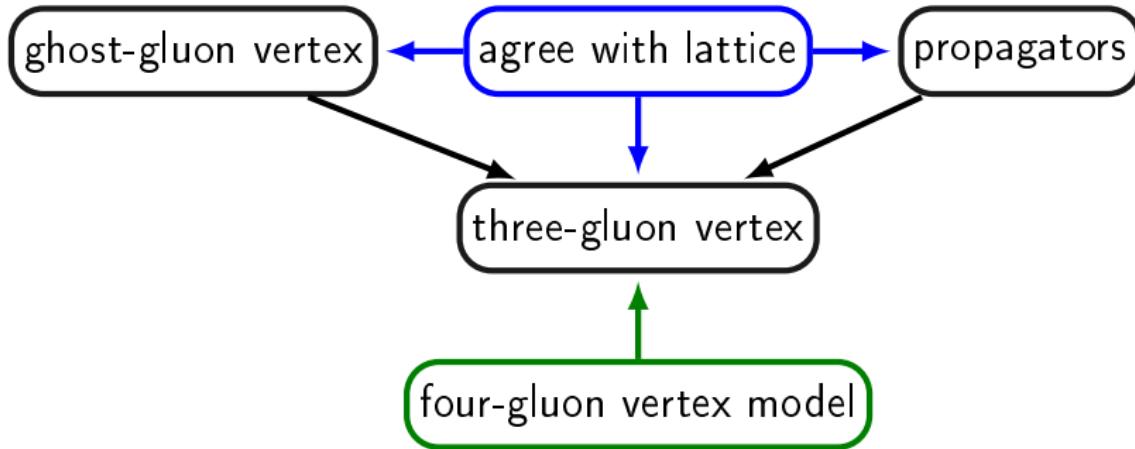


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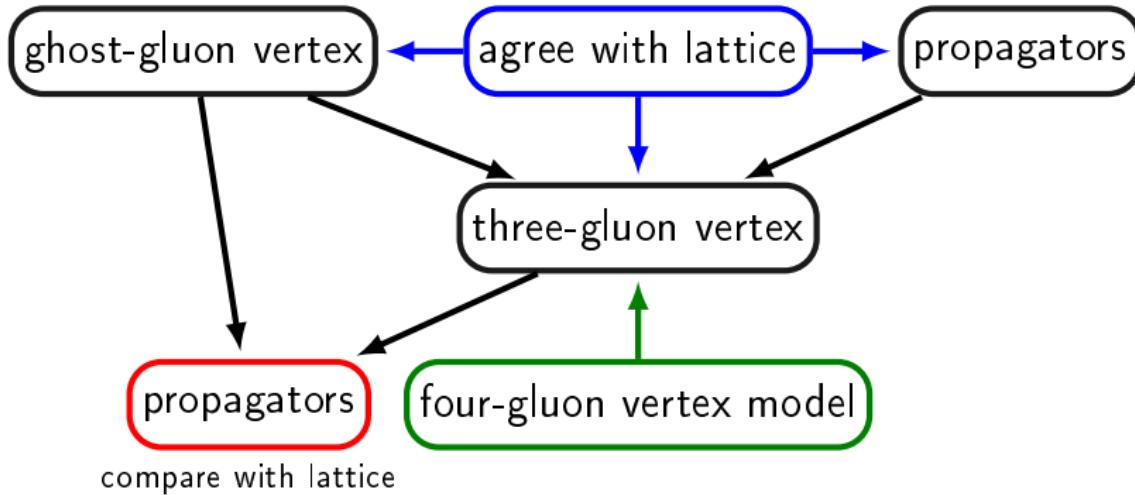
Strategy



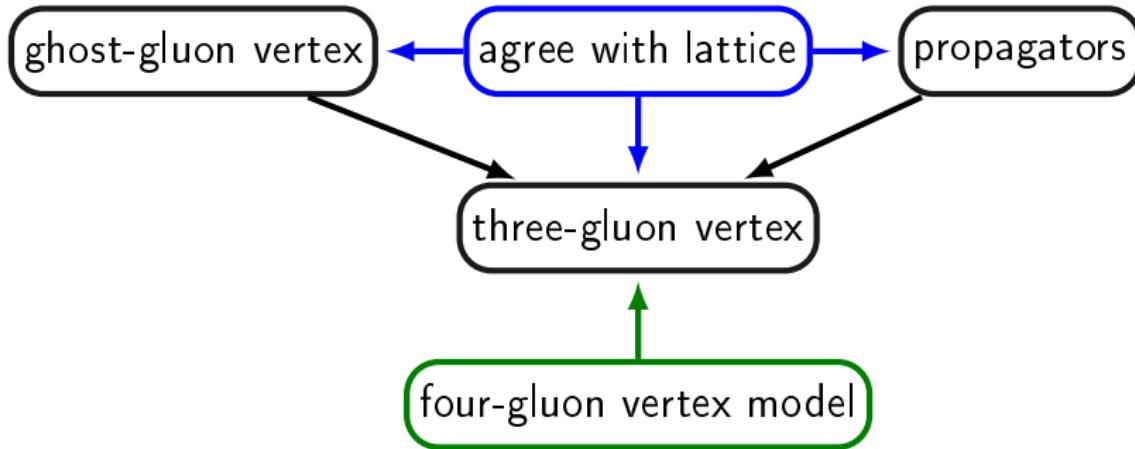
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Strategy



The three-gluon vertex |

Four-gluon vertex model (decoupling):

$$D^{A^4}(p, q, r, s) = (a \tanh(b/\bar{p}^2) + 1) D_{RG}^{A^4}(p, q, r, s)$$

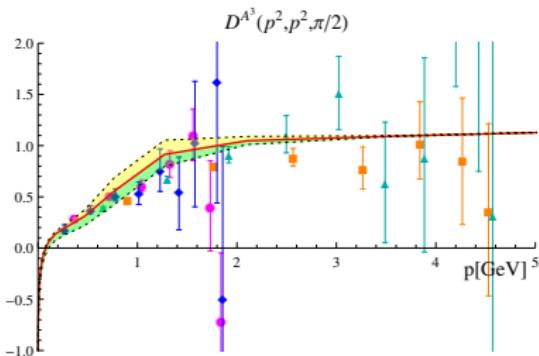
→ Test model dependence by varying a and b .

The three-gluon vertex I

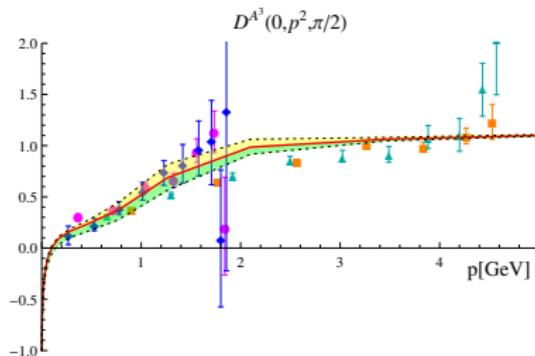
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[Blum, MQH, Mitter, von Smekal '14; lattice: Cucchieri, Maas, Mendes '08]



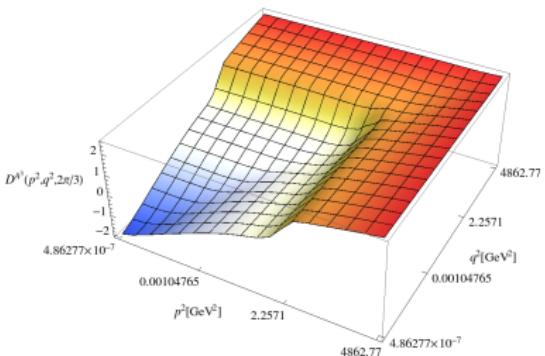
→ Truncation reliable. Neglected terms, including two-loop, suppressed.

See also results by [Eichmann, Williams, Alkofer, Vujinovic '14], esp. other dressings, and [Peláez, Tissier, Wschebor '13].

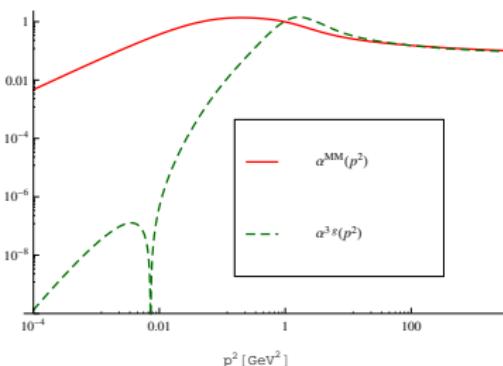
The three-gluon vertex II

$$\Gamma_{\mu\nu\rho}^{AAA,abc}(p, q, k) := i g f^{abc} D^{AAA}(p^2, q^2, \cos\theta) \Gamma_{\mu\nu\rho}^{AAA,(0)}(p, q, k)$$

Fixed angle:

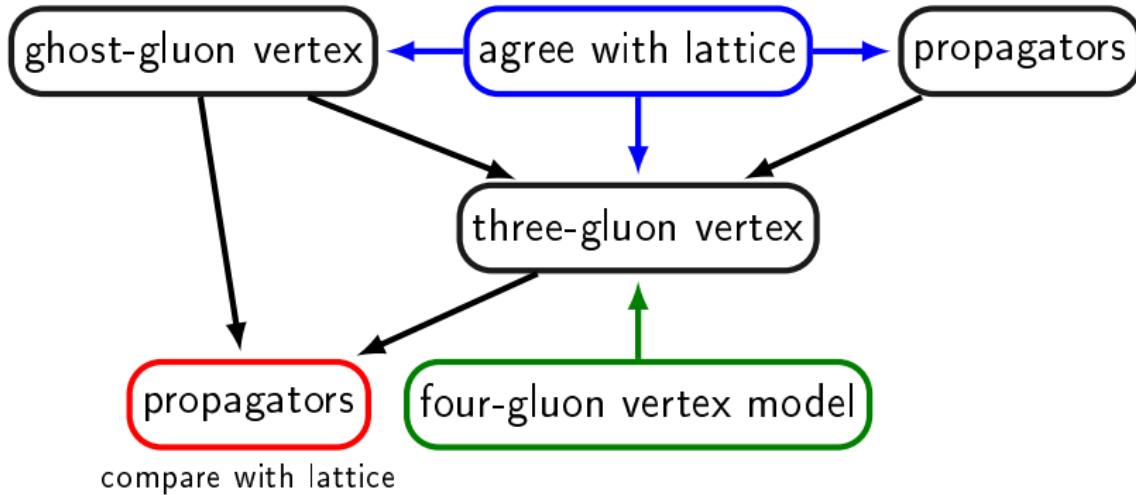


Couplings:

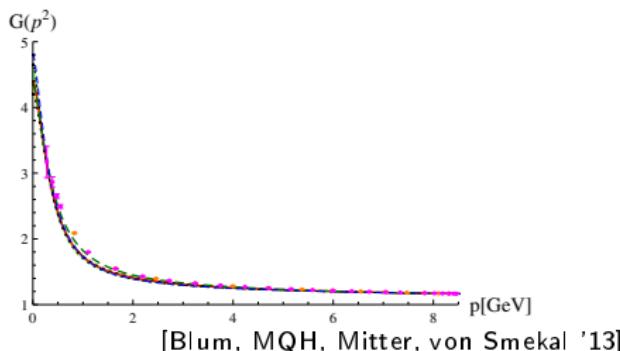
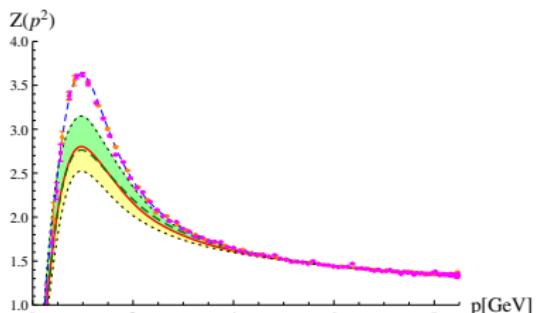


- Bose symmetry visible (enforced by hand).
- Zero crossing (in this tensor), also found by [Peláez, Tissier, Wschebor '13; Aguilar, Binosi, Ibáñez, Papavassiliou '13; Eichmann, Williams, Alkofer, Vujinovic '14] and on lattice in $d = 2, 3$.

Strategy



Propagators II: Limits of one-loop truncation



[Blum, MQH, Mitter, von Smekal '13]

- Ghost almost unaffected.
- Gap in midmomentum regime must be due to missing two-loop diagrams!

NB: Employed projection of three-gluon vertex is the same as in gluon loop of gluon propagator DSE! \rightarrow Error from neglected tensors negligible.

Explicit two-loop studies [Bloch '03; Mader, Alkofer '12; Meyers, Swanson '14]:
squint \gg sunset diagram

Four-gluon vertex

- 6 external variables
- 4 integration variables

Cf. propagator: 1 ext., 2 int.

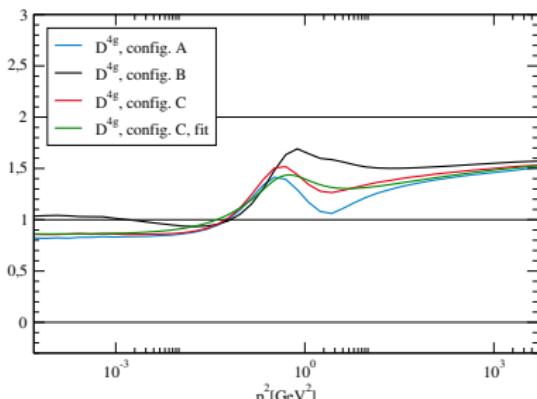
2 propagators → laptop

four-gluon vertex → > 100 cores on cluster

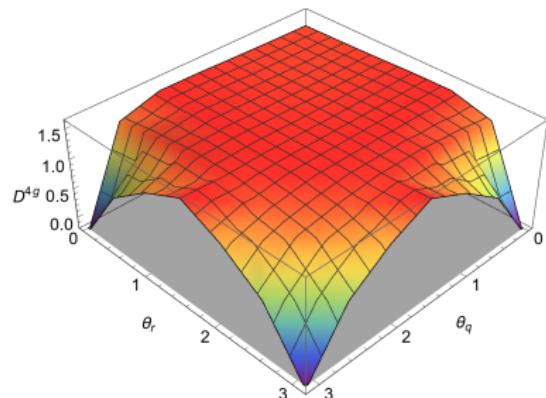
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[Cyrol, MQH, von Smekal '14]



2-parameter fit:

$$D_{\text{model}}^{4g, \text{ dec}}(p, q, r, s) = (\text{atanh}(b/\bar{p}^2) + 1) D_{\text{RG}}^{4g}(p, q, r, s)$$

Beyond Landau gauge: Coulomb gauge

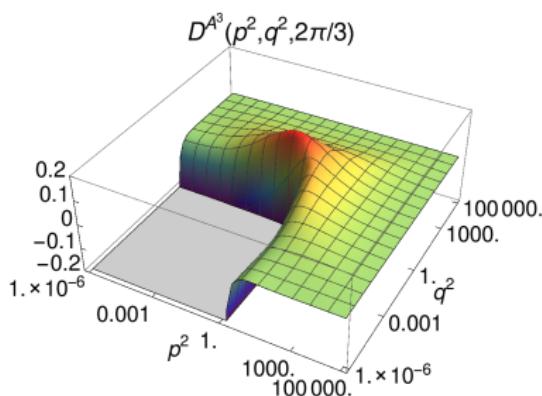
Why the Landau gauge is convenient

- Minimum number of terms in DSEs.
- Transversality → longitudinal part decouples.
- Historically ghost-gluon vertex provided the entry point (special here).

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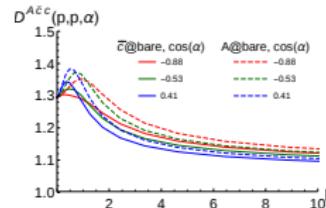
[MQH, Campagnari, Reinhardt '14]

Three-gluon vertex:

- Zero crossing
- IR divergent like p^{-3}

Ghost-gluon vertex:

- Different truncations quite similar



Summary

Highlighted prospects:

- Description of **hadron properties** from first principles.
- **QCD phase diagram**: explore non-zero density.

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- 2-, 3- and 4-point functions calculated
- Truncation effects understood (after more than 30 years!)
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- System of DSEs closes with this truncation
→ **self-contained, quantitative description.**

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Thank you for your attention.