

Infrared Behavior of the Three-Gluon Vertex in d-Dimensional Yang-Mills Theory

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Motivation

Diploma thesis: IR-exponents of vertices in three dimensions

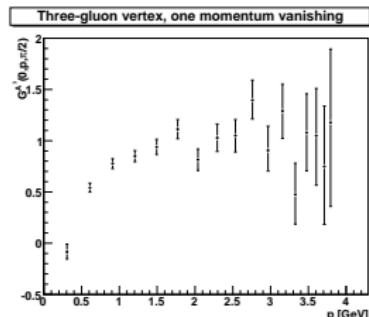
Why d space-time dimensions?

	4 dimensions	3 dimensions
Ordering scheme	skeleton expansion Alkofer, Fischer, Llanes-Estrada	ordering scheme?
Lattice	hardly feasible in IR	Cucchieri, Maas, Mendes

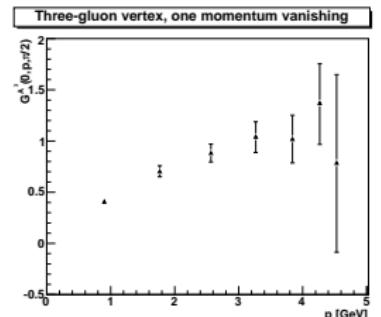
Later?:

- Three-dimensional Yang-Mills (YM) theory can be compared to the infinite-temperature limit of four-dimensional YM theory.
- A connection between 4d YM in Coulomb gauge and 3d YM in Landau gauge exists.

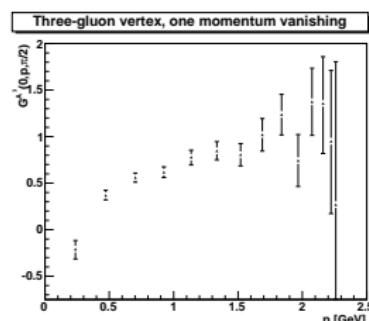
Available Lattice Data



Maas, in preparation



Cucchieri et al., hep-lat/0610006



Cucchieri et al., PRD74, 014503, 2006

The infrared region is easier to access in two and three dimensions.

What is an infrared exponent?

Infrared exponent expresses the behavior of a Green function at small external momenta ($\ll \Lambda_{QCD}$).

Green functions can have additional momentum dependence,

e.g. ghost propagator: $D(p^2) = -\frac{G(p^2)}{p^2}$

Prominent picture in QCD: Gribov-Zwanziger scenario

- gluon propagator dressing function vanishes in IR: $Z(p^2) \propto (p^2)^{2\kappa}$
- ghost propagator dressing function is divergent in IR: $G(p^2) \propto (p^2)^{-\kappa}$
- $\kappa_{d=4} = 0.595$

Outline

Transversality of gluon propagator



Bare ghost-gluon vertex



Infrared exponents of the propagators from the ghost Dyson-Schwinger equation (DSE)



Skeleton expansion of other DSEs



Infrared exponents of all dressed Green functions of Yang-Mills theory

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The Bare Ghost-Gluon Vertex

Starting point is transversal gluon propagator in Landau gauge:

$$k_\mu D_{\mu\nu}(k) = k_\mu \frac{Z(k)}{k^2} \left[\delta_{\mu\nu} - \frac{k_\mu k_\nu}{k^2} \right] = 0$$

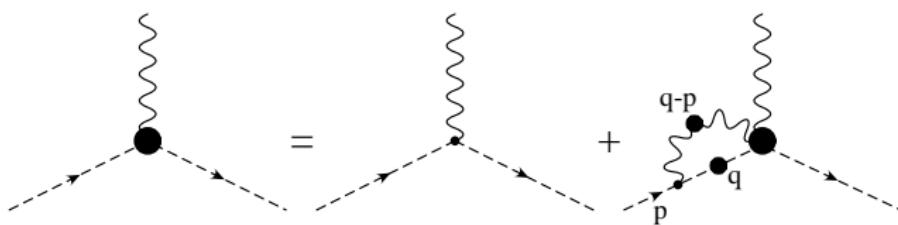
$$(q-p)_\mu D_{\mu\nu}(q-p) = 0 \Rightarrow q_\mu D_{\mu\nu}(q-p) = p_\mu D_{\mu\nu}(q-p)$$

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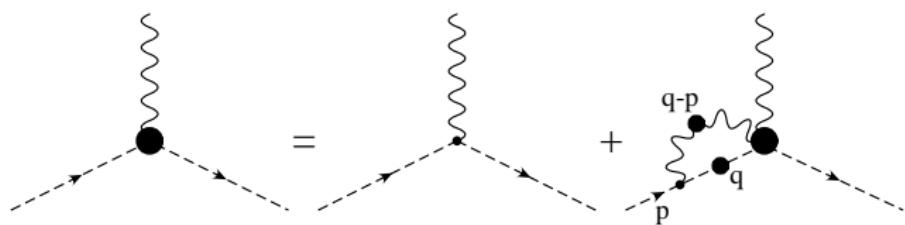
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loop momentum q and external momentum p

The Bare Ghost-Gluon Vertex

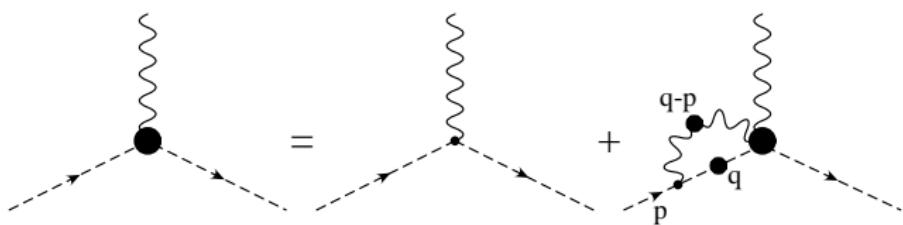
$$q_\mu D_{\mu\nu}(q - p)$$



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The Bare Ghost-Gluon Vertex

$$q_\mu D_{\mu\nu}(q - p) \rightarrow 0 \text{ for } p \rightarrow 0$$



loop momentum q and external momentum p

The Bare Ghost-Gluon Vertex

$$q_\mu D_{\mu\nu}(q - p) \rightarrow 0 \text{ for } p \rightarrow 0$$

- Consequence 1: Ghost-gluon vertex stays bare in the IR.
- Consequence 2: Its renormalization constant \tilde{Z}_1 can be set to 1.

Definition of renormalization constant:

$$G_\mu(q, p)|_{p_1^2 = p_2^2 = p_3^2 = 0} = \tilde{Z}_1 G_\mu^{(bare)}(q, p).$$

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The Propagators

Power law ansatz for dressing functions:

$$Z(p^2) = A \cdot (p^2)^\alpha$$

$$G(p^2) = B \cdot (p^2)^\beta$$



Comparison of IR exponent of lhs and rhs:

$$1 - \beta = \frac{d}{2} + \alpha - 1 + \beta - 1 + \frac{1}{2} + \frac{1}{2}$$

$$\alpha = -2\beta + 2 - \frac{d}{2}$$

Set $\beta = -\kappa$: $\rho_{2,1} = 0$, $\rho_{2,0} = -\kappa$, $\rho_{0,2} = 2\kappa + 2 - \frac{d}{2}$

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n-Ghost-n-Antighost-m-Gluon Vertex

l number of loops

m_i number of internal gluons

n_i number of internal ghosts

$v_{0,3}$ number of dressed three-gluon vertices

$v_{0,3}^b$ number of bare three-gluon vertices

$v_{2,1}$ number of ghost-gluon vertices, dressed or bare

$v_{0,4}^b$ number of bare four-gluon vertices

v momentum dimension of the bare n-point function itself, e.g.
three-gluon vertex: $\frac{1}{2}$

$$\begin{aligned}\rho_v = & (l - m_i + v_{0,3}) \frac{d}{2} + (2m_i - n_i - 3v_{0,3})\kappa + \\ & + \frac{1}{2}(2m_i - 2n_i - 3v_{0,3} + v_{0,3}^b + v_{2,1} - 2v)\end{aligned}$$

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Higher Orders and Other Diagrams

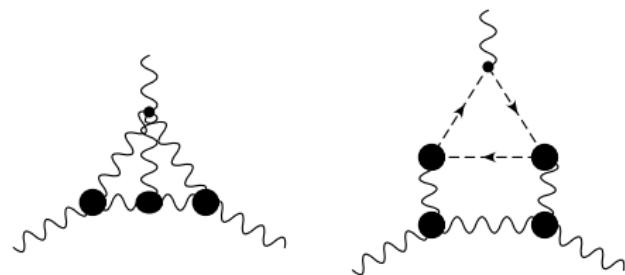
The general formula still contains bare vertices:

$$\begin{aligned}\rho_{n,m} = & (-n + 1 - v_{0,4}^b - v_{0,3}^b) \frac{d}{2} + (4v_{0,4}^b + 3v_{0,3}^b - m + n)\kappa \\ & + (2v_{0,4}^b + 2v_{0,3}^b + 2n - 2)\end{aligned}$$

Contributions of these,

$$v_{0,3}^b \left(3\kappa - \frac{d}{2} + 2 \right)$$

$$v_{0,4}^b \left(4\kappa - \frac{d}{2} + 2 \right),$$



are always positive \Rightarrow diagrams not dominant

$$\rho_{n,m} = (n - m)\kappa + (1 - n)\left(\frac{d}{2} - 2\right)$$

d=4: Alkofer et al., Phys. Lett. B611, 2005
Fischer 2006, unpublished

Differences Between the Dimensions

IR exponent:

$$\rho_{n,m} = (-m+n)\kappa + (1-n)\left(\frac{d}{2} - 2\right)$$

The full scaling:

Dimension	d	4	3	2
κ		0.6 ¹	0.4/0.5 ²	0.2/0 ²
Ghost	$-\kappa - 1$	-1.6	-1.4/1.5	-1.2/1
Gluon	$2\kappa + 1 - \frac{d}{2}$	0.2	0.3/0.5	0.4/0
3-gluon	$-3\kappa + \frac{d}{2} - \frac{3}{2}$	-1.3	-1.2/1.5	-1.1/0.5
4-gluon	$-4\kappa + \frac{d}{2} - 2$	-2.4	-2.1/2.5	-1.8/1

⇒ Qualitative behavior is the same in two, three and four dimensions.

¹Lerche 2001; Zwanziger, Phys. Rev. D6, 2002

²Maas et al., Eur. Phys. J. C37, 2004; Zwanziger, Phys. Rev. D6, 2002

Conclusions

- The skeleton expansion is valid in all dimensions.
- The qualitative behavior in the infrared is the same in two, three and four dimensions, namely dominated by ghost contributions.