Glueballs from functional equations



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Institute of Theoretical Physics Giessen University MQH, Phys.Rev.D 101, <u>arXiv:2003.13703</u> MQH, Fischer, Sanchis-Alepuz, Eur.Phys.J.C 80, <u>arXiv:2004.00415</u>

A Virtual Tribute to *Quark Confinement and the Hadron Spectrum* 2021, virtually in Stavanger, Norway, Aug. 6, 2021



QCD bound states

Bound states in QCD



Mesons

Baryons

QCD bound states

Bound states in QCD



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Glueball observations

Experimental candidates, but situation not conclusive.

Scalar glueball: 0^{++} , mixing with scalar isoscalar mesons

Candidate reaction: $J/\psi \rightarrow \gamma + 2g$

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Recent analysis of BESIII data [Sarantsev, Denisenko, Thoma, Klempt '21]:



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Glueballs

Glueball calculations

Yang-Mills theory

- "Isolated" problem: only gluons
- Clean picture: well-established lattice results

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QCD glueballs: mixing with quarks

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QCD glueballs: mixing with quarks

Unquenching on the lattice [Gregory et al. '12]:

- Much higher statistics required (poor signal-to-noise ratio)
- Continuum extrapolation and inclusion of fermionic operators still to be done
- Mixing with q
 q
 q
 challenging
- Tiny (e.g., 0^{++} , 2^{++}) to moderate unquenching effects (e.g., 0^{-+}) found
- $m_{\pi}=360\,\mathrm{MeV}$

Glueball calculations

Yang-Mills theory

- "Isolated" problem: only gluons
- Clean picture: well-established lattice results
- Functional methods: High quality input available for bound state equations

QCD glueballs: mixing with quarks

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Hadrons from bound state equations



Integral equation: $\Gamma(q, P) = \int dk \, \Gamma(k, P) \, S(k_{+}) \, S(k_{-}) \, K(k, q, P)$

Hadrons from bound state equations





Bound state equations

Glueball BSE



Need \ldots and solve for \rightarrow . \rightarrow Mass

Bound state equations

Glueball BSE



Need (000) and $(1, solve for) \rightarrow Mass$ Not quite...

Glueball BSE



Gluons couple to ghosts \rightarrow Include 'ghostball'-part. (First step: no quarks \rightarrow Yang-Mills theory)

Glueball BSE



Need $(\mathfrak{M}, \rightarrow)$ and $4\times$, solve for \rightarrow and \rightarrow . \rightarrow Mass

Construction of kernel

Consistency with input: Apply same construction principle.

Glueball BSE



Need $(\Omega, --$ and $4\times$, solve for \rightarrow and \rightarrow . \rightarrow Mass

Construction of kernel

Consistency with input: Apply same construction principle.

Previous BSE calculations for glueballs:

- [Meyers, Swanson '13]
- [Sanchis-Alepuz, Fischer, Kellermann, von Smekal '15]
- [Souza et al. '20]
- [Kaptari, Kämpfer '20]

⇒ Input is important for quantitative predictive power!

[MQH, Fischer, Sanchis-Alepuz '20]

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Bound state equations

Kernel construction

From 3PI effective action truncated to three-loops:

[Fukuda '87; McKay, Munczek '89; Sanchis-Alepuz, Williams '15; MQH, Fischer, Sanchis-Alepuz '20]



- Some diagrams vanish for certain quantum numbers.
- Full QCD: Same for quarks \rightarrow Mixing with mesons.

Equations

Equations of motion from 3-loop 3PI effective action



Gluon and ghost fields: Elementary fields of Yang-Mills theory in the Landau gauge

Self-contained system of equations with the scale as the only input.

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Truncation \rightarrow 3-loop expansion of 3PI effective action [Berges '04]

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Gluon and ghost fields: Elementary fields of Yang-Mills theory in the Landau gauge

Self-contained system of equations with the scale as the only input.

Truncation \rightarrow 3-loop expansion of 3PI effective action [Berges '04]

- 4 coupled integral equations with full kinematic dependence.
- Sufficient numerical accuracy required for renormalization.
- One- and two-loop diagrams [Meyers, Swanson '14; MQH '17; Eichmann, Pawlowski, Silva '21].

Results

Landau gauge propagators

Gluon dressing function:



- Family of solutions: Nonperturbative completions of Landau gauge [Maas '10]?
- Realized by condition on G(0) [Fischer, Maas, Pawlowski '08; Alkofer, MQH, Schwenzer '08]
- Results here independent of G(0)

Gluon propagator:



Ghost dressing function:



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Results

Concurrence of functional methods

Exemplified with three-gluon vertex.

3PI vs. 2-loop DSE:



[Cucchieri, Maas, Mendes '08; Sternbeck et al. '17; MQH '20]

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DSE vs. FRG:



[Cucchieri, Maas, Mendes '08; Sternbeck et al. '17; Cyrol et al. '16; MQH '20]

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DSF vs. FRG:

Concurrence of functional methods

Exemplified with three-gluon vertex.

3PI vs. 2-loop DSE:



Beyond this truncation

- Further dressings of three-gluon vertex [Eichmann, Williams, Alkofer, Vujinovic '14]
- Effects of four-point functions [MQH '16, MQH '17, Corell et al. '18, MQH '18]

Jeballs

BSE

Solving a BSE



ueballs

BSE

Solving a BSE



Consider the eigenvalue problem (Γ is the BSE amplitude)

 $\mathcal{K} \cdot \Gamma(\mathcal{P}) = \lambda(\mathcal{P}) \Gamma(\mathcal{P}).$

 $\lambda(P^2) = 1$ is a solution to the BSE \Rightarrow Glueball mass $P^2 = -M^2$

ueballs

BSE

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Calculation requires quantities for

$$k_{\pm}^2 = P^2 + k^2 \pm 2\sqrt{P^2 k^2} \cos \theta = -M^2 + k^2 \pm 2 i M \sqrt{k^2} \cos \theta.$$

 \Rightarrow Complex momentum arguments.

Direct calculation from functional methods possible, e.g., [Fischer, MQH '20].

 \rightarrow talk by Windisch

Alternative

Extrapolate λ from $P^2 > 0$.

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Extrapolation of \lambda(P^2)
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Extrapolation method

- Extrapolation to time-like P² using Schlessinger's continued fraction method (proven superior to default Padé approximants) [Schlessinger '68]
- Average over extrapolations using subsets of points for error estimate

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Coefficients a_i can determined such that f(x) exact at x_i .

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Test extrapolation for solvable system: Heavy meson [MQH, Sanchis-Alepuz, Fischer '20]





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Method

Extrapolation of $\lambda(P^2)$ for glueballs

Higher eigenvalues: Excited states.

lueballs

Method

Extrapolation of $\lambda(P^2)$ for glueballs

Higher eigenvalues: Excited states.



Physical solutions for $\lambda(P^2) = 1$.

Glueballs

Results

Glueballs masses for $0^{\pm+}$



All results for $r_0 = 1/418(5)$ MeV.

[MQH, Fischer, Sanchis-Alepuz '20]

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alueballs Results

Two-loop diagrams

Results from [MQH, Fischer, Sanchis-Alepuz '20] were from one-loop terms only:





Fully self-consistent DSE/BSE truncation

 \rightarrow two-loop terms (complete 3-loop truncated 3PI effective action)

lueballs Results

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Drastic increase in computational resources, hence lower precision used.

Preliminary result for 0^{++} , 0^{-+} : No effect on mass.

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Results

Glueball masses for $J^{\pm+}$

For higher spin, larger tensor bases: more tensors, more indices

Glueballs Re

Results

Glueball masses for $J^{\pm +}$

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[MQH, Fischer, Sanchis-Alepuz, in preparation]

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• Quantitatively reliable correlation functions (Euclidean) from functional equations



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 - Comparison with lattice results



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Thank you for your attention.

More details...

Glueballs as bound states

Hadron masses from correlation functions of color singlet operators.

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Hadron masses from correlation functions of color singlet operators. Example: For $J^{PC} = 0^{++}$ glueball take $O(x) = F_{\mu\nu}(x)F^{\mu\nu}(x)$:

 $D(x - y) = \langle O(x)O(y) \rangle$

- $\bullet \rightarrow$ Lattice: Mass from this correlator by exponential Euclidean time decay.
- Complicated object in a diagrammatic language: 2-, 3- and 4-gluon contributions

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Put total momentum on-shell and consider individual 2-, 3- and 4-gluon contributions. \rightarrow Each can have a pole at the glueball mass.

 A^4 -part of D(x - y), total momentum on-shell:



More details...

Charge parity

Transformation of gluon field under charge conjugation:

$$A^a_\mu
ightarrow -\eta(a) A^a_\mu$$

where

$$\eta(a) = \begin{cases} +1 & a = 1, 3, 4, 6, 8\\ -1 & a = 2, 5, 7 \end{cases}$$

Color neutral operator with two gluon fields:

$$A^a_\mu A^a_
u o \eta(a)^2 A^a_\mu A^a_
u = A^a_\mu A^a_
u.$$

 $\Rightarrow C = +1$

Negative charge parity, e.g.:

$$egin{aligned} d^{abc} A^a_\mu A^b_
u A^c_
ho &
ightarrow - d^{abc} \eta(a) \eta(b) \eta(c) A^a_\mu A^b_
u A^c_
ho &= \ - d^{abc} A^a_\mu A^b_
u A^c_
ho. \end{aligned}$$

Only nonvanishing elements of the symmetric structure constant d^{abc}: zero or two indices equal to 2, 5 or 7.

More details...

Landau gauge propagators in the complex plane

Simpler truncation:



More details..

Landau gauge propagators in the complex plane

Simpler truncation:





 \rightarrow Opening at $q^2 = p^2$.

More details..

Landau gauge propagators in the complex plane



Appearance of branch cuts for complex momenta forbids integration directly to cutoff.

Deformation of integration contour necessary [Maris '95]. Recent resurgence: [Alkofer et al. '04; Windisch, MQH, Alkofer, '13; Williams '19; Miramontes, Sanchis-Alepuz '19; Eichmann et al. '19, ...]

More details.

Landau gauge propagators in the complex plane

Ray technique for self-consistent solution of a DSE:



- Current truncation leads to a pole-like structure in the gluon propagator.
- Analyticity up to 'pole' confirmed by various tests (Cauchy-Riemann, Schlessinger, reconstruction)
- No proof of existence of complex conjugate poles due to simple truncation.

[Fischer, MQH '20]

More details. .

Landau gauge vertices





- Nontrivial kinematic dependence of ghost-gluon vertex
- Simple kinematic dependence of three-gluon vertex
- Four-gluon vertex from solution

Four-gluon vertex:

Three-gluon vertex:

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[MQH '20]

More details..

Some properties of the Landau gauge solution

 Slavnov-Taylor identities (gauge invariance): Vertex couplings agree down to GeV regime

[MQH '20]

More details..

Some properties of the Landau gauge solution

 Slavnov-Taylor identities (gauge invariance): Vertex couplings agree down to GeV regime

 Renormalization: First parameter-free subtraction of quadratic divergences
 ⇒ One unique free parameter (family of solutions)

